



# El problema inverso en la astrofísica de hoyos negros

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# Outline



Description of an inverse problem

Various examples

Brutal force vs intelligence methods

Inverse / Direct problems involving BHs

Binary black holes from the strain

Kicked black hole

Wandering black hole

Intrinsic parameters out of an image

# Direct and Inverse problem



Consider a phenomenon described by

$$\text{given } x \text{ define } F(x) =: y$$

Here  $x$  may represent initial conditions or physical parameters.

And  $F$  represents the model or theory assumed to rule the phenomenon.

Direct problem. Given  $x$  and  $F$  calculate  $y$ . Easy  $y = F(x)$ .

Inverse problem of the cause. Given  $F$  and  $y$ , find  $x$  that produces  $y$ . It could be  $x = F^{-1}(y)$

Inverse problem of the model. Given  $y$  and values  $x$ , find  $F$  such that  $F(x)$  produces  $y$ .

Inverse problem of the cause and the model. Given  $y$ , find  $F$  and  $x$  such that  $F(x)$  produces  $y$ .

# Examples of direct problem



Given  $m$ ,  $b$ ,  $k$ ,  $x(0)$  and  $\dot{x}(0)$  determine  $x(t)$  when it obeys

$$m\ddot{x} + b\dot{x} + kx = 0$$

Given the initial position  $\mathbf{x}(0)$  of a particle of mass  $m$  moving around an object of mass  $M$ , calculate  $\mathbf{x}(t)$  for  $t > 0$ , provided

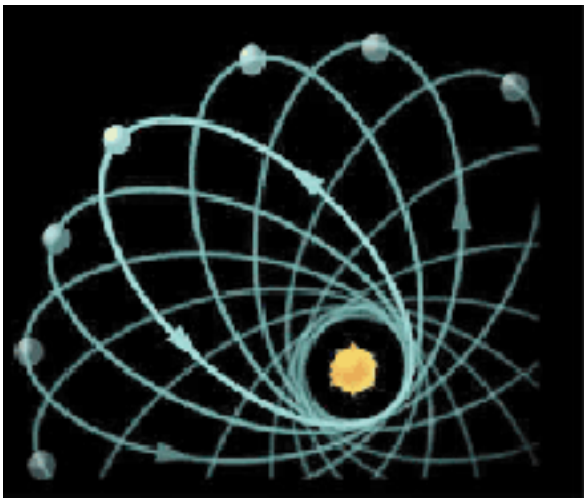
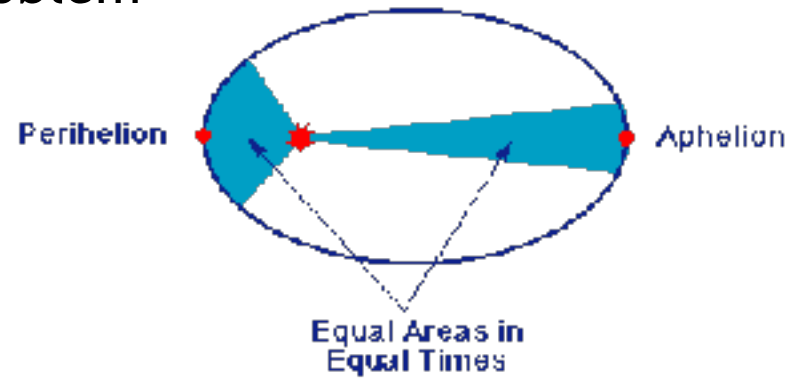
$$\mathbf{F} = G \frac{Mm}{|\mathbf{x}|^3} \mathbf{x} = m\ddot{\mathbf{x}}$$

*These are simple PVI's*



# Examples of Inverse problems / data fitting

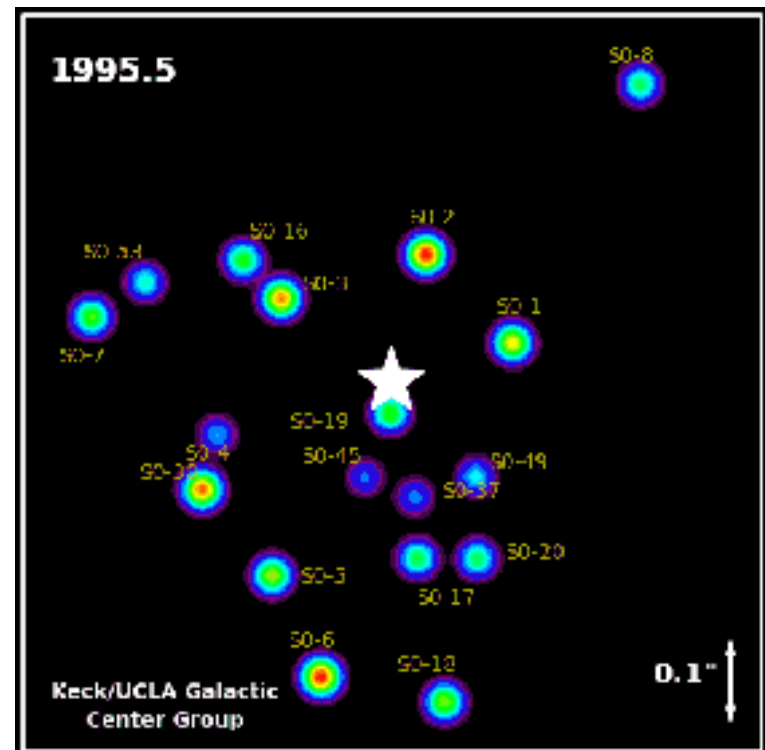
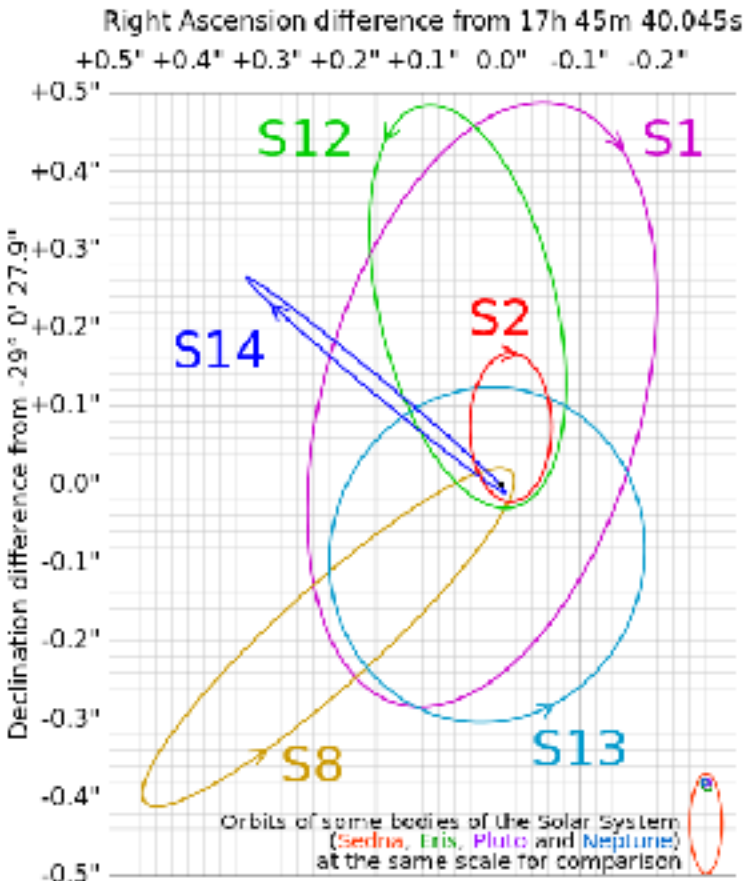
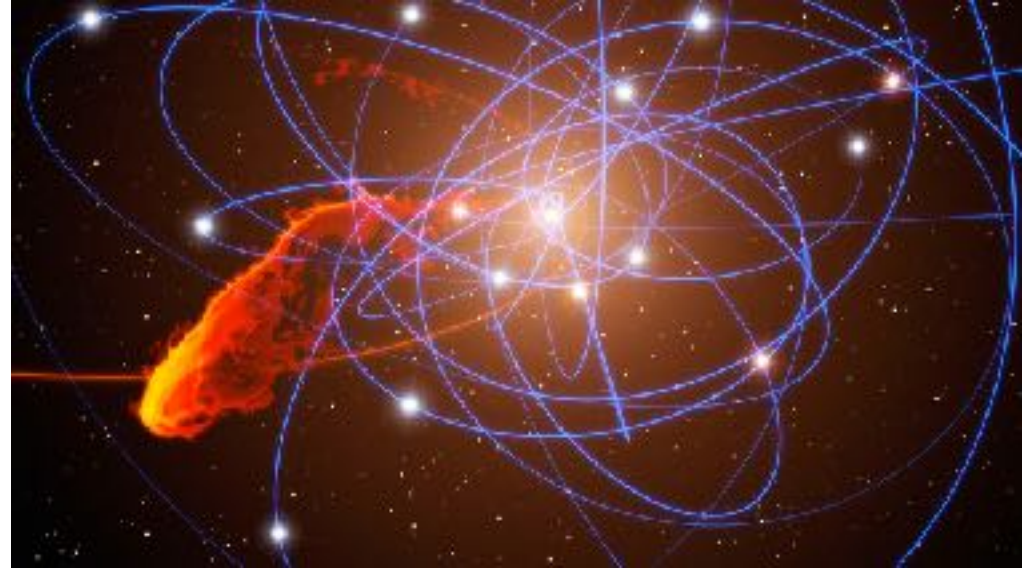
## Kepler's problem / Newton problem



Inverse problem

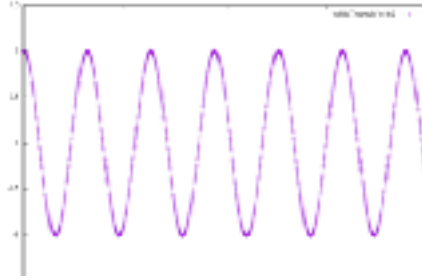
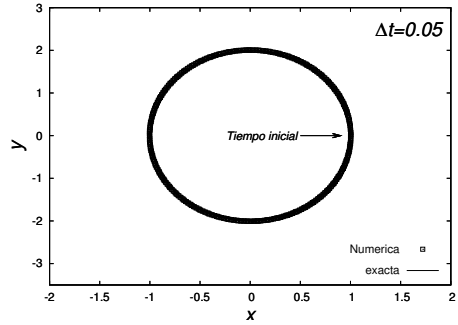
1. Of the cause provided Newton's law
2. Of the model as it happened to be

# Example: Sgr A\*

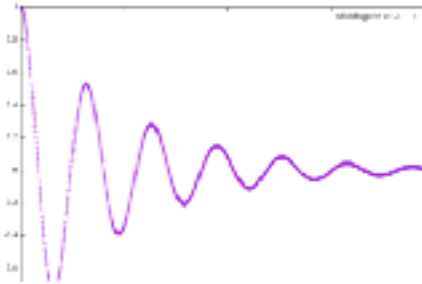
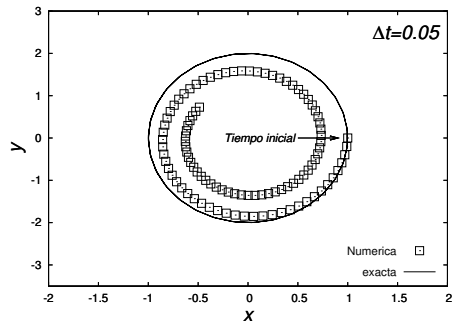




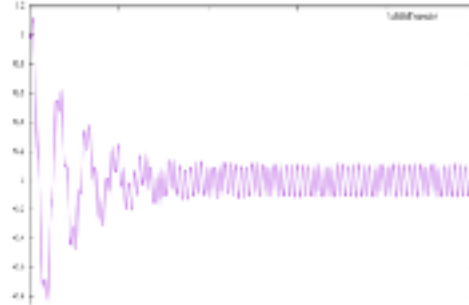
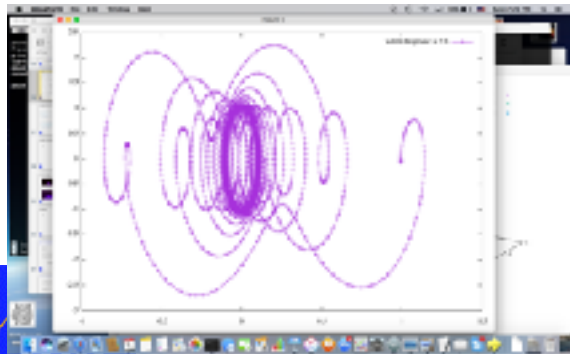
# Finding a law and initial conditions: linear problems



$$m\ddot{x} = -kx$$

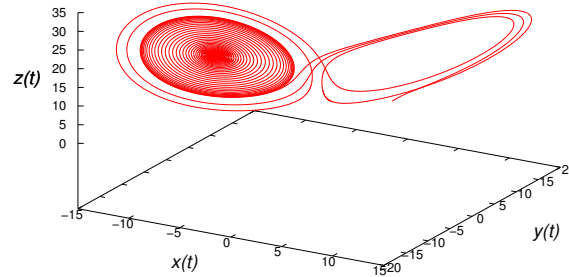
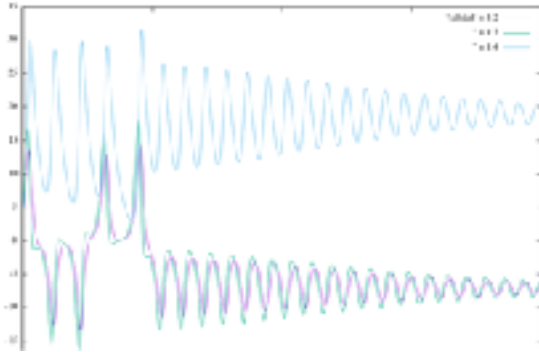


$$\ddot{x} + \frac{b}{m}\dot{x} + \frac{k}{m}x = 0.$$

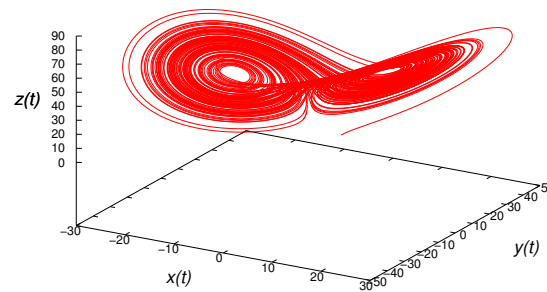
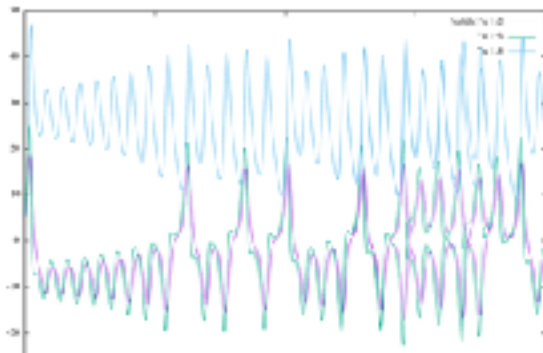




Solving Inverse problems depend A LOT on  $F$  (when there is one): see non-linear cases



$$\begin{aligned} \frac{dx}{dt} &= a(y - x) \\ \frac{dy}{dt} &= rx - y - xz \\ \frac{dz}{dt} &= xy - bz \end{aligned}$$

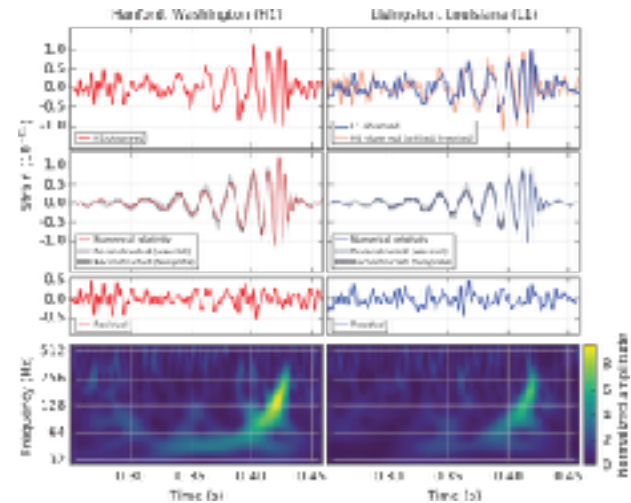
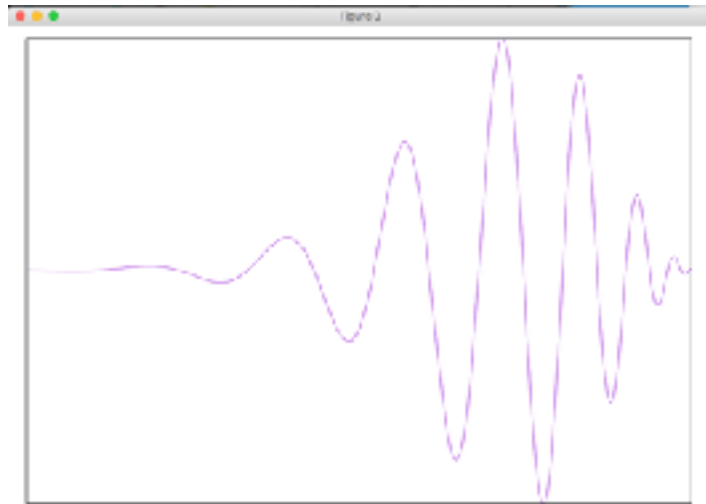


$$\begin{aligned} a &= 10 & d_t &= 0.001 \\ b &= 8/3 & x_0 = y_0 = z_0 &= 5.0 \\ r &= 10, 15, 50 & 0 \leq t \leq 100 & \end{aligned}$$

**Imagine when the system is ruled by PDEs: computer power, time, time, time**



# Modeling data vs data science



```

subroutine colcchs(my_t,my_u)
  use numbers
  implicit none
  real(kind=8), intent(in) :: my_t
  real(kind=8), dimension(NE), intent(in) :: my_u

  rhs(1) = my_u(2)
  rhs(2) = - b / m * my_u(2) - k / m * my_u(1) + F0 / m * sin( CapOmega * my_t ) * exp(-(my_t-50.)**2/20.**2)
end subroutine colcchs
    
```

$$\begin{aligned}
 ds^2 &= -\alpha^2 dt^2 + \gamma_a(dt^2 + \beta^a dx^a + \beta^a dx^a) \\
 \phi &= \frac{1}{2} \ln \gamma & \gamma_a &= e^{-2\sigma} \gamma_a \\
 K &= \gamma_a K^a & A_a &= e^{-2\sigma} (K_a - \frac{1}{2} \gamma_a K^b) \\
 \tilde{t}^a &= \gamma^a \tilde{t}_a = -\beta^a \tilde{t}^a
 \end{aligned}$$

$$\begin{aligned}
 (\partial_t - \mathcal{L}_\beta) \gamma_a &= -2\sigma_a \dot{\sigma} \\
 (\partial_t - \mathcal{L}_\beta) \sigma &= -\frac{1}{2} \sigma K \\
 (\partial_t - \mathcal{L}_\beta) \dot{\sigma}_a &= e^{-2\sigma} (-D_t \dot{\sigma}_a - \alpha \partial_t \gamma^b (\dot{\sigma}_b - \sigma \delta_b \dot{\sigma}_a - 2\dot{\sigma}_b \dot{\sigma}_a^b)) \\
 (\partial_t - \mathcal{L}_\beta) K &= -D^a \dot{\sigma}_a + \alpha (\dot{\sigma}_a \dot{\sigma}^a - \frac{1}{2} K^2) \\
 \partial_t^2 &= 2\alpha (\tilde{t}_a \dot{\sigma}^a - \alpha \tilde{t}^a \partial_t \sigma - \frac{1}{2} \gamma^a \partial_t K) - 2\tilde{t}^a \partial_t \sigma + \gamma^a \partial_t \dot{\sigma}_a \\
 &\quad + \frac{1}{2} \gamma^a \partial_t \sigma_a \dot{\sigma}^a + \sigma \partial_t \tilde{t}^a + \frac{1}{2} \gamma^a \partial_t \dot{\sigma}^a - \frac{-(\gamma^a + \frac{1}{2} (\tilde{t}^a - \sigma \gamma^a)) \alpha \dot{\sigma}^a}{\gamma_a \gamma_b \gamma_c}
 \end{aligned}$$

Yoneda et al. (2002)



# Solving Inverse problems depends A LOT on $F$ (when there is one)

## On the capacity to solve direct problems

And  $F$  is really challenging as one approaches a realistic model

Systems ruled by ODEs are actually friendly, one can still run millions of combinations of parameters and initial conditions

Then estimate one of the combinations with the minimum error

$$\begin{aligned}
 ds^2 &= -\alpha^2 dt^2 + \gamma_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt) \\
 \phi &= \frac{1}{2} \ln \gamma & \gamma_{ij} &= e^{-2\phi} \gamma_{ij} \\
 K &= \gamma_{ij} K^{ij} & A_{ij} &= e^{-2\phi} (K_{ij} - \frac{1}{2} \gamma_{ij} K) \\
 \dot{\Gamma}^i &= \gamma^{ij} \dot{\Gamma}_{jk} = -\partial_t \gamma^{ij}
 \end{aligned}$$

$$\begin{aligned}
 (\partial_t - \mathcal{L}_\beta) \gamma_{ij} &= -2\alpha A_{ij} \\
 (\partial_t - \mathcal{L}_\beta) \phi &= -\frac{1}{2} \alpha K \\
 (\partial_t - \mathcal{L}_\beta) \dot{\Gamma}_{ij} &= e^{-2\phi} [ -D_i D_j \phi - \alpha \partial_{ij} \Gamma^k - \alpha (K A_{ij} - 2A_{ik} A_{jl}) \\
 (\partial_t - \mathcal{L}_\beta) K &= -D^i D_i \alpha + \alpha (A_{ij} A^{ij} - \frac{1}{2} K^2) \\
 \partial_t \Gamma^i &= 2\alpha ( \dot{\Gamma}_{jk} A^{ij} - \alpha \dot{\Gamma}^i \partial_j \phi - \frac{1}{2} \gamma^{ij} \partial_j K ) - 2A^{ik} \partial_j \alpha + \gamma^{ij} \partial_j \partial_k \phi \\
 &\quad + \frac{1}{2} \gamma^{ij} \partial_k \alpha \partial_j \phi + \beta^k \partial_{ij} \Gamma^i + \frac{1}{2} \gamma^{ij} \partial_j \alpha \partial_k \phi \quad \text{--- (Yoneda et al., 2002) ---}
 \end{aligned}$$

$$\frac{\partial \rho}{\partial t} + \nabla \cdot (\rho \mathbf{v}) = 0,$$

$$\frac{\partial (\rho \mathbf{v})}{\partial t} + \nabla \cdot \left( \left( p + \frac{1}{2} \mathbf{B}^2 \right) \mathbf{I} + \rho \mathbf{v} \mathbf{v} - \mathbf{B} \mathbf{B} \right) = -(\nabla \cdot \mathbf{B}) \mathbf{B} + \rho \mathbf{g},$$

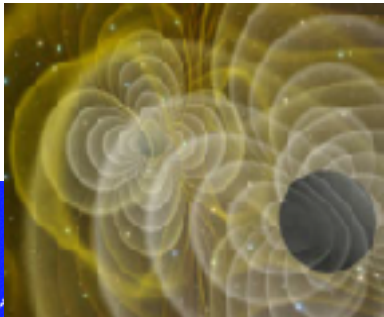
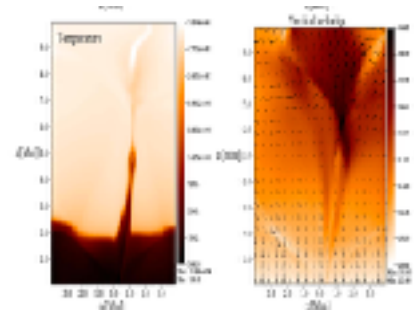
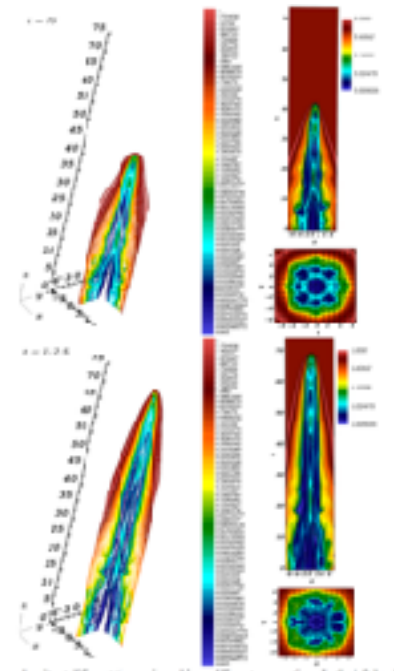
$$\frac{\partial E}{\partial t} + \nabla \cdot \left( \mathbf{v} \left( E + \frac{1}{2} \mathbf{B}^2 + p \right) - \mathbf{B} (\mathbf{B} \cdot \mathbf{v}) \right) = -\mathbf{B} \cdot (\nabla \psi) - \nabla \cdot ((\eta \cdot \mathbf{J}) \times \mathbf{B}) + \rho \mathbf{g} \cdot \mathbf{v},$$

$$\frac{\partial \mathbf{B}}{\partial t} + \nabla \cdot (\mathbf{B} \mathbf{v} - \mathbf{v} \mathbf{B} + \psi \mathbf{I}) = -\nabla \times (\eta \cdot \mathbf{J}),$$

$$\frac{\partial \psi}{\partial t} + c_h^2 \nabla \cdot \mathbf{B} = -\frac{c_h^2}{c_p^2} \psi,$$

$$\mathbf{J} = \nabla \times \mathbf{B},$$

$$E = \frac{p}{(\gamma - 1)} + \frac{\rho v^2}{2} + \frac{\mathbf{B}^2}{2},$$



# Methods of solution



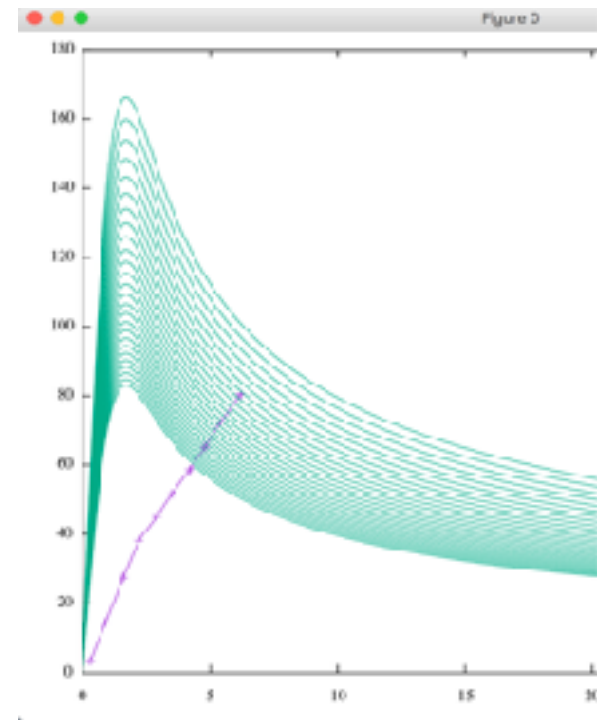
Brutal force:

Works fine for system ruled by ODEs (for instance)

Does NOT work for complicated PDEs

Genetic algorithms

Artificial Intelligence methods

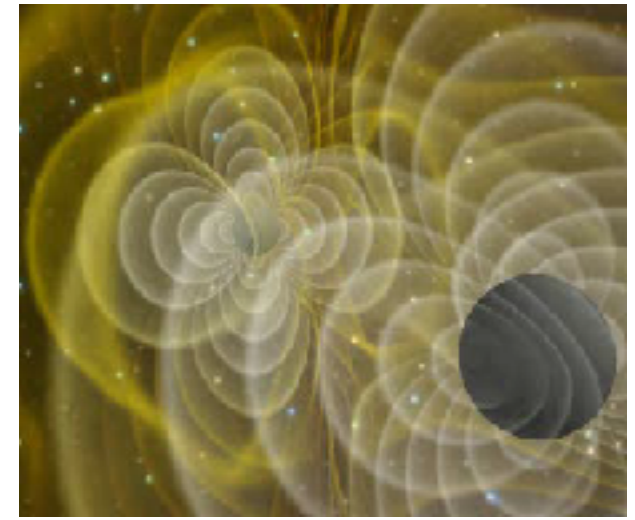
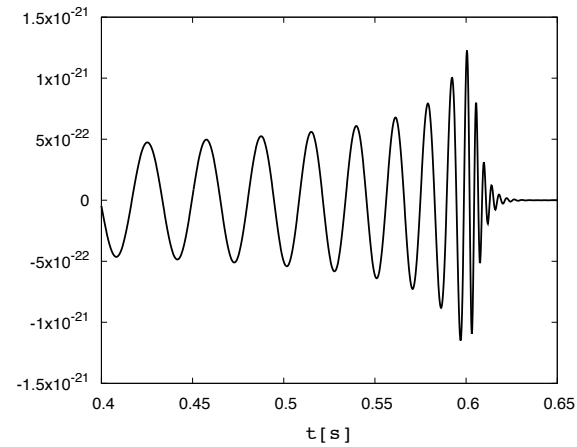
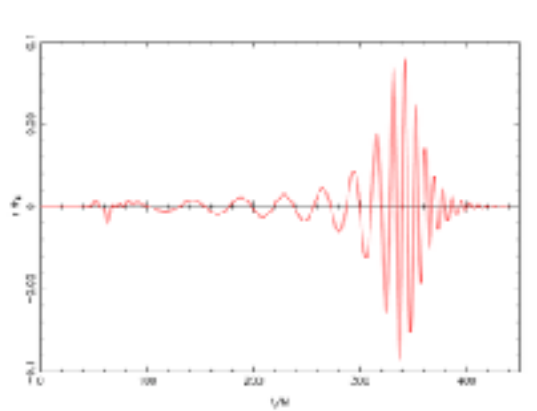




# Inverse Problem 1: BBH ... the direct problem first

In a simulation one estimates the Weyl Psi4 scalar

It is decomposed in spherical harmonics and get the dominant modes



$$h = h_+ - ih_x = \int_{-\infty}^t dt' \int_{-\infty}^{t'} dt'' \Psi_4$$

$$ds^2 = -\alpha^2 dt^2 + \gamma_{ij}(dx^i + \beta^i dt)(dx^j + \beta^j dt)$$

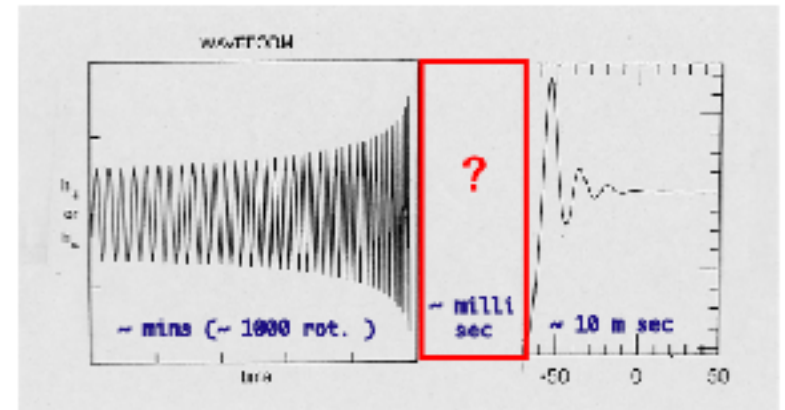
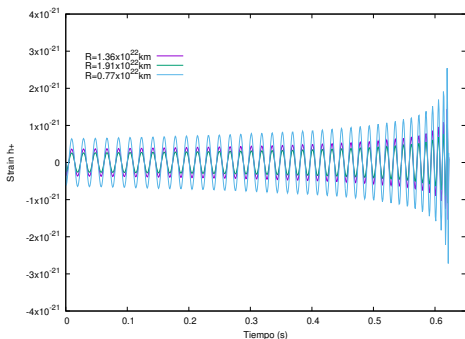
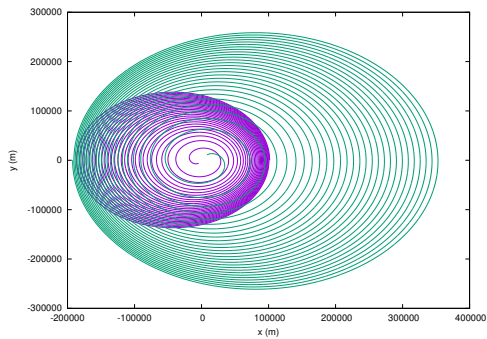
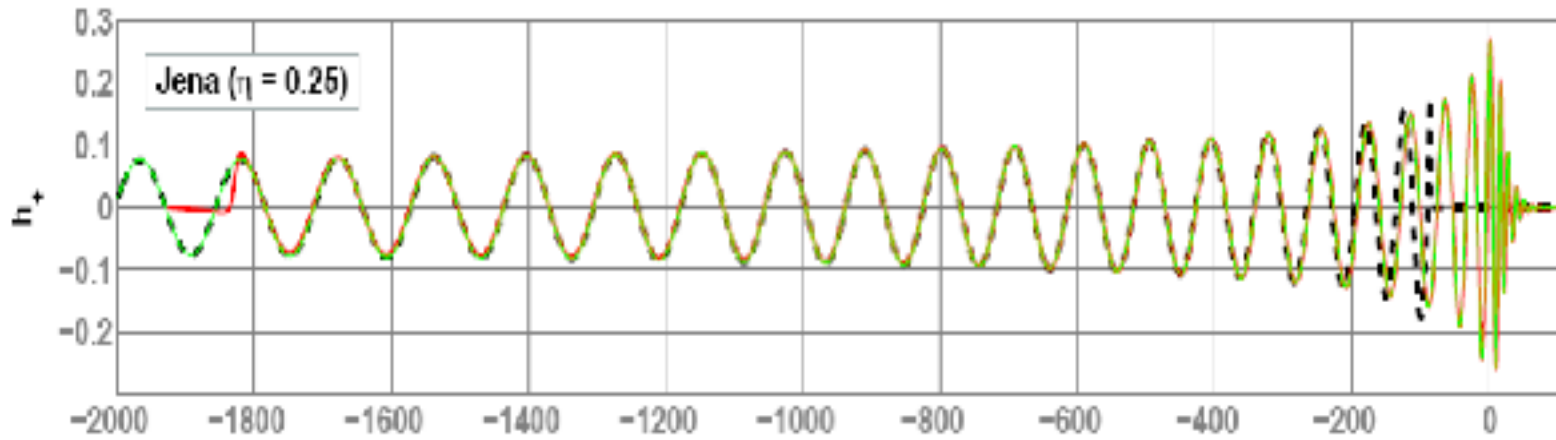
$$\begin{aligned} \phi &= \frac{1}{2}(\gamma + \gamma_1) & \gamma_1 &= e^{-4u} \gamma_0 \\ K &= \gamma_0 K^0 & \tilde{A}_0 &= e^{-4u} (K_0 - \frac{1}{2} \gamma_0 K) \\ \tilde{\Gamma} &= \gamma^0 \tilde{\Gamma}_0 & &= -\partial_t \tilde{\Gamma}^0 \end{aligned}$$

$$\begin{aligned} (\partial_t - \mathcal{L}_\beta) \gamma_0 &= -2\alpha \tilde{A}_0 \\ (\partial_t - \mathcal{L}_\beta) \beta &= -\frac{1}{2} \alpha K \\ (\partial_t - \mathcal{L}_\beta) \tilde{A}_0 &= e^{-2u} (-D_0 D_0 \alpha - \alpha \partial_0 \tilde{\Gamma}^0 - \alpha) K \tilde{A}_0 - 2 \tilde{A}_0 \partial_0^2 \alpha \\ (\partial_t - \mathcal{L}_\beta) K &= -D^0 D_0 \alpha + \alpha (\tilde{A}_0 A^0 - \frac{1}{2} K^2) \\ \partial_t^2 &= 2\alpha \left( \tilde{\Gamma}_0^0 \tilde{A}^0 - \partial_0^2 \gamma_0 - \frac{1}{2} \gamma_0 \partial_0 K \right) - 2 \tilde{A}_0 \partial_0 \alpha + \alpha^2 \partial_0 \partial_0 \alpha \\ &\quad + \frac{1}{2} \alpha^2 \partial_0 \partial_0 K + \alpha \partial_0 \tilde{\Gamma}^0 + \frac{1}{2} \tilde{\Gamma}_0^0 \partial_0^2 \alpha \end{aligned}$$

To be published (2002)



# Inverse Problem 1: BBH ... the direct problem first



**INSPIRAL** → **COALESCE** → **BLACKHOLE FORMATION**  
 Innermost Stable  
 Circular Orbit?

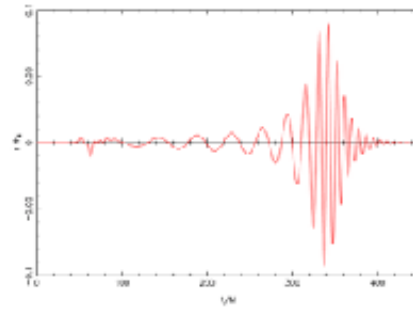
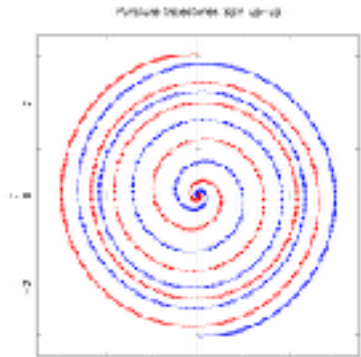
Post Newtonian Approx.

↳ Numerical Relativity

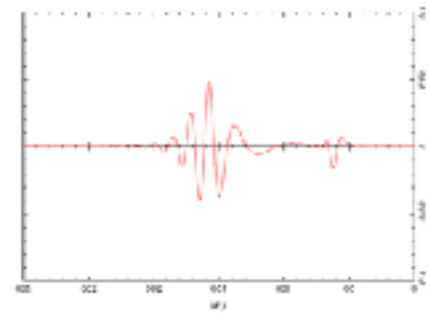
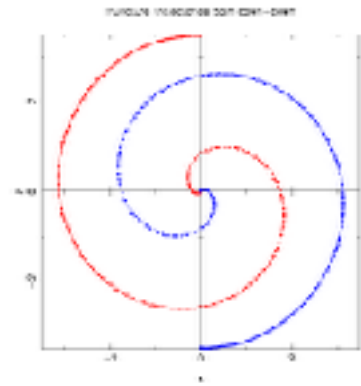
↳ BH. Perturbation



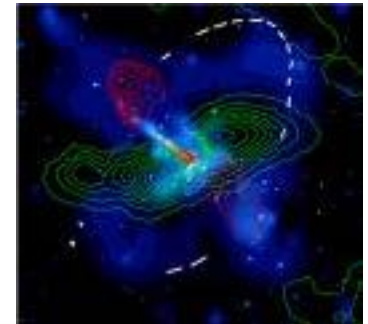
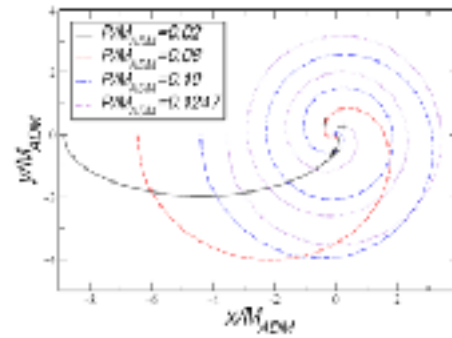
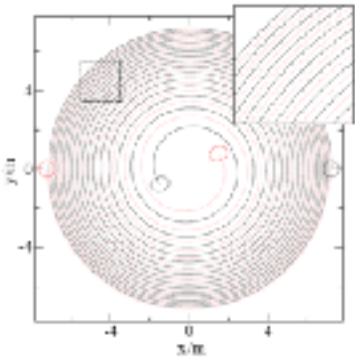
# Direct problem: exploration a nightmare



$\uparrow s_1 \quad \uparrow L \quad \uparrow s_2$



$\downarrow s_1 \quad \uparrow L \quad \downarrow s_2$



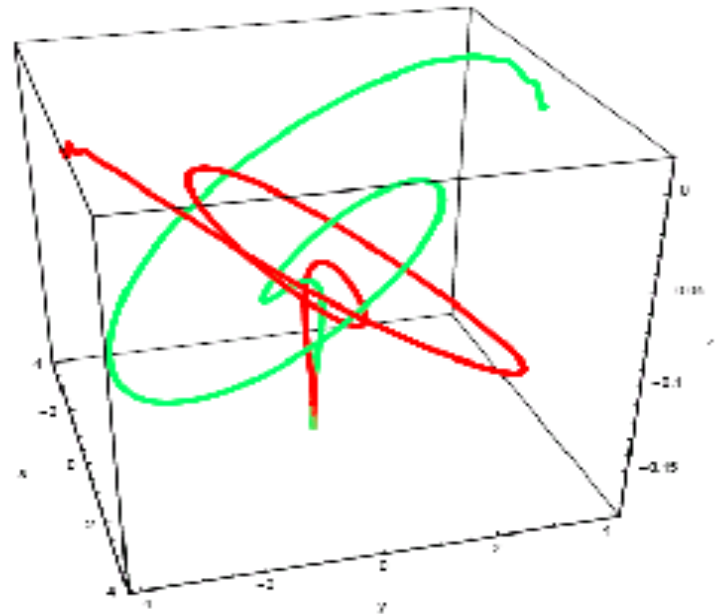
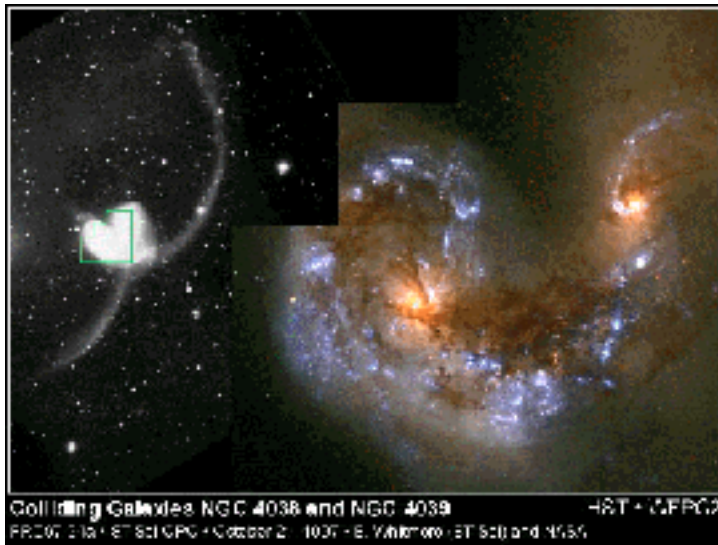
[guzman@ifm.umich.mx](mailto:guzman@ifm.umich.mx)



# Direct problem: kicks

Escape velocities:

Globulares	clusters	30Km/s
dSph		20-100km/2
dE		100-300km/s
giant galacies		1000km/2





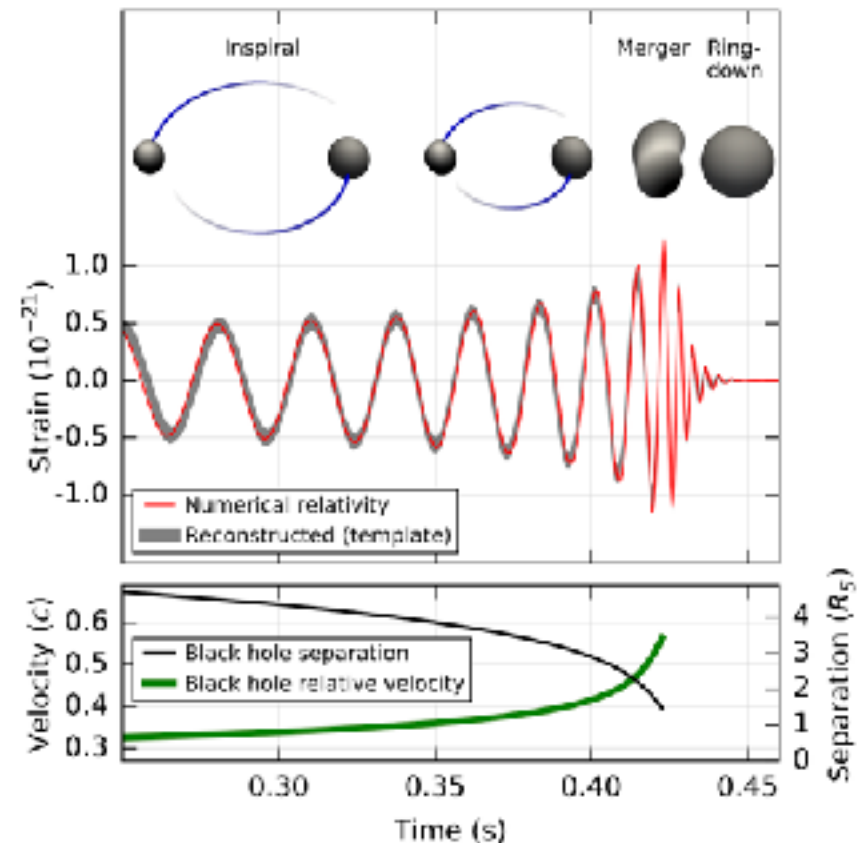
# Numerical Relativity: TOO SLOW - too expensive

PostNewtonian approximation: *leading order*

$$h^{insp}(t) = \frac{4GM\eta}{D_L c^2} \left( \frac{GM}{c^3} \frac{d\phi}{dt} \right)^{2/3} \cos[2\phi(t)]$$

$D_L$  es la distancia a la fuente

$$\eta = \mu / M$$





# Plus ringdown

## Ringdown

$$h(t) = \sum_{n=0}^{N-1} A_n e^{-i\sigma_n(t-t_{match})}$$

$n$  overtone

$N$  number of overtones considered

$A_n$  Amplitudes determining the MATCHING

$$\sigma_n = \omega_n - i\alpha_n$$

$$\omega_n > 0$$

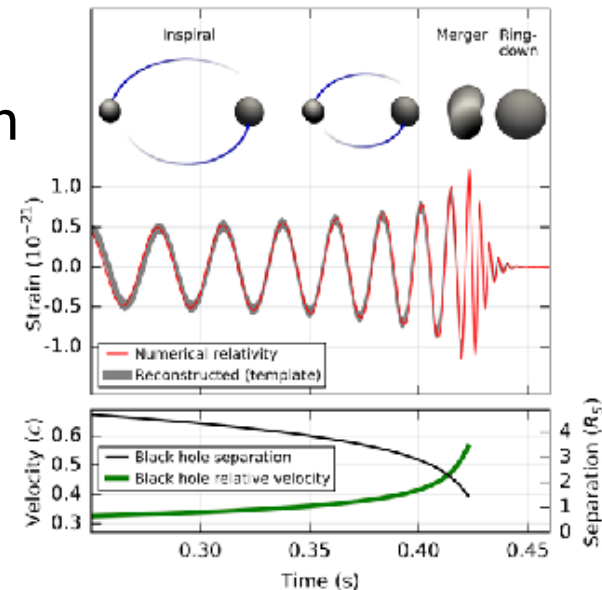
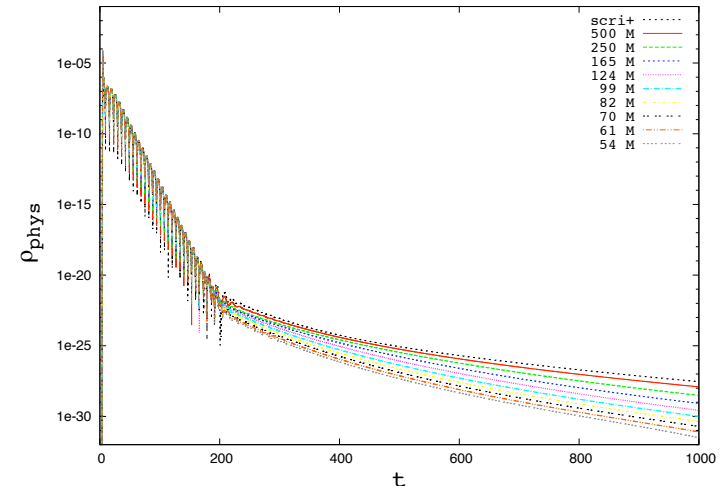
son las frecuencias de oscilación

$$\alpha_n > 0$$

son el inverso del tiempo de decaimien

These quantities depend on the MASS

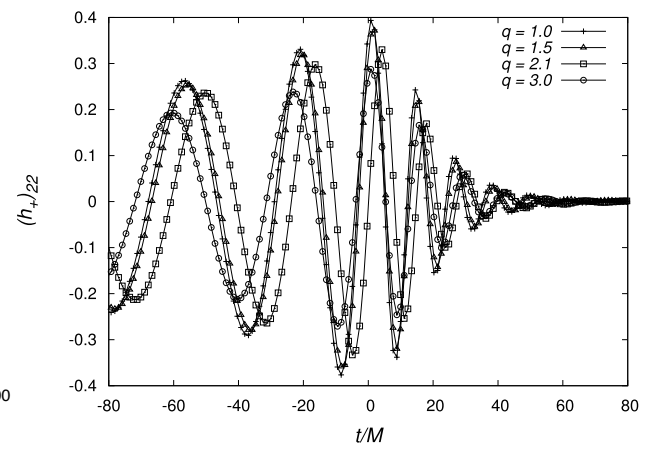
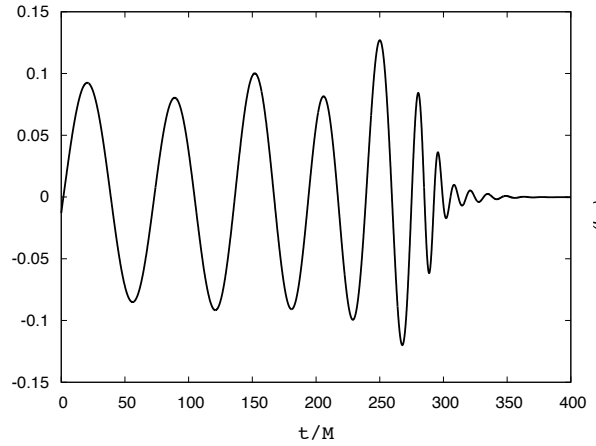
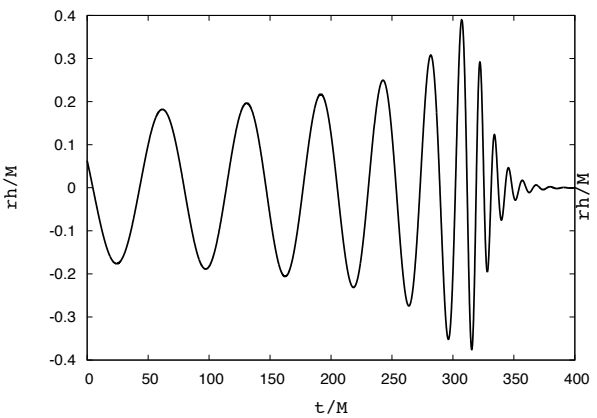
And SPIN of the **FINAL BLACK HOLE**





# Inverse Problem 1: BBH parameters from GW strain

$$M = m_1 + m_2, q = \frac{m_1}{m_2}, \chi_1 = \frac{S_1}{m_1^2}, \chi_2 = \frac{S_2}{m_2^2}$$



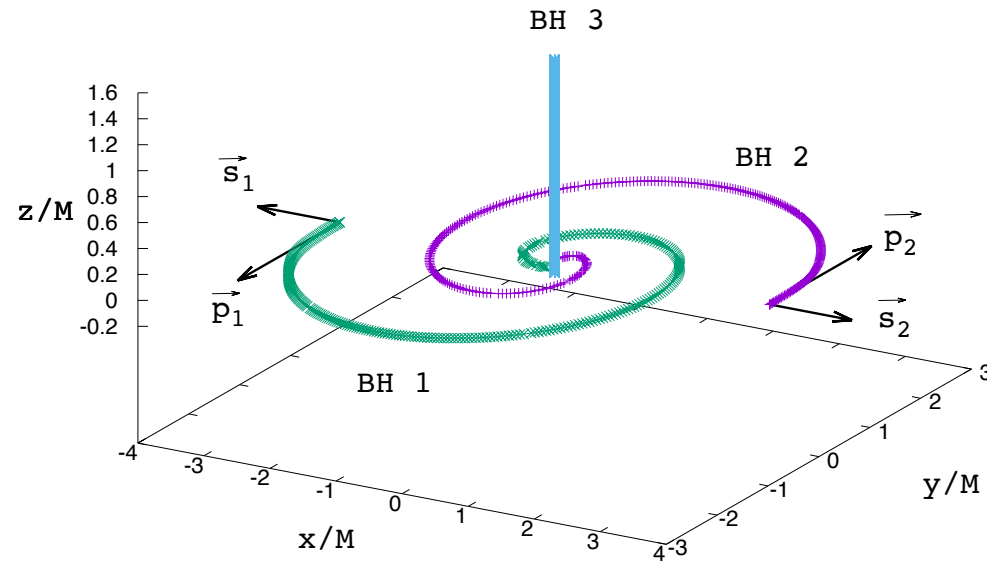
Parameter estimates in binary black hole collisions using neural networks  
M. Carrillo, M. Gracia-Linares, J. A. González, F. S. Guzmán  
Gen. Rel. Grav. 48, 141 (2016)

One parameter Binary Black Hole inverse problem using a sparse training set  
M. Carrillo, M. Gracia-Linares, J. A. González, F. S. Guzmán  
Int. J. Mod. Phys. D 27, 1850043 (2018)



# Inverse Problem 2: kicked BH inverse problem

Of the cause  
Of the cause + model



	0.15	0.18
(1)	$p_{1y} = p_{1z} = 0.133$	$v = \frac{0.0047 \cdot 1.4776 \cdot 10^8 \text{ km}}{4.927 \text{ s}} = 2008.5 \text{ km/s}$
(2)	$p_{1y} = p_{1z} = 0.120$	$v = -0.001544 ( ) = 462.855 \text{ km/s}$
(3)	$p_{1y} = p_{1z} = 0.125$	$v = 0.00354248 (299776.74) = 1061 \text{ km/s}$
(A)	$p_{1y} = p_{1z} = 0.124$	$v = 0.001913 (299776.74) = 573.47 \text{ km/s}$
(5)	$p_{1y} = p_{1z} = 0.123$	$v = 0.006$ AGAIN!!!
(6)	$p_{1y} = p_{1z} = 0.124$	$v = 0.0056!!!$
(7)	$p_{1y} = p_{1z} = 0.122$	$v = 0.00487$
(8)	$p_{1y} = p_{1z} = 0.1215$	$v = 0.00372 = 1115.17 \text{ km/s}$
(9)	$p_{1y} = p_{1z} = 0.1212$	$v = 0.002518 = 759.85 \text{ km/s}$
(7)	$p_{1y} = p_{1z} = 0.1218$	$v = 0.00318 = 953.3$

$$\vec{V}_{\text{recoil}}(q, \vec{\alpha}) = v_m \hat{e}_1 + v_{\perp} (\cos \xi \hat{e}_1 + \sin \xi \hat{e}_2) + v_{\parallel} \hat{n}_{\parallel},$$

$$v_m = A \frac{\eta^2 (1-q)}{(1+q)} [1 + B \eta],$$

$$v_{\perp} = H \frac{\eta^2}{(1+q)} \left[ (1 + B_H \eta) (\alpha_2^{\parallel} - q \alpha_1^{\parallel}) + H_S \frac{(1-q)}{(1+q)^2} (\alpha_2^{\parallel} + q^2 \alpha_1^{\parallel}) \right],$$

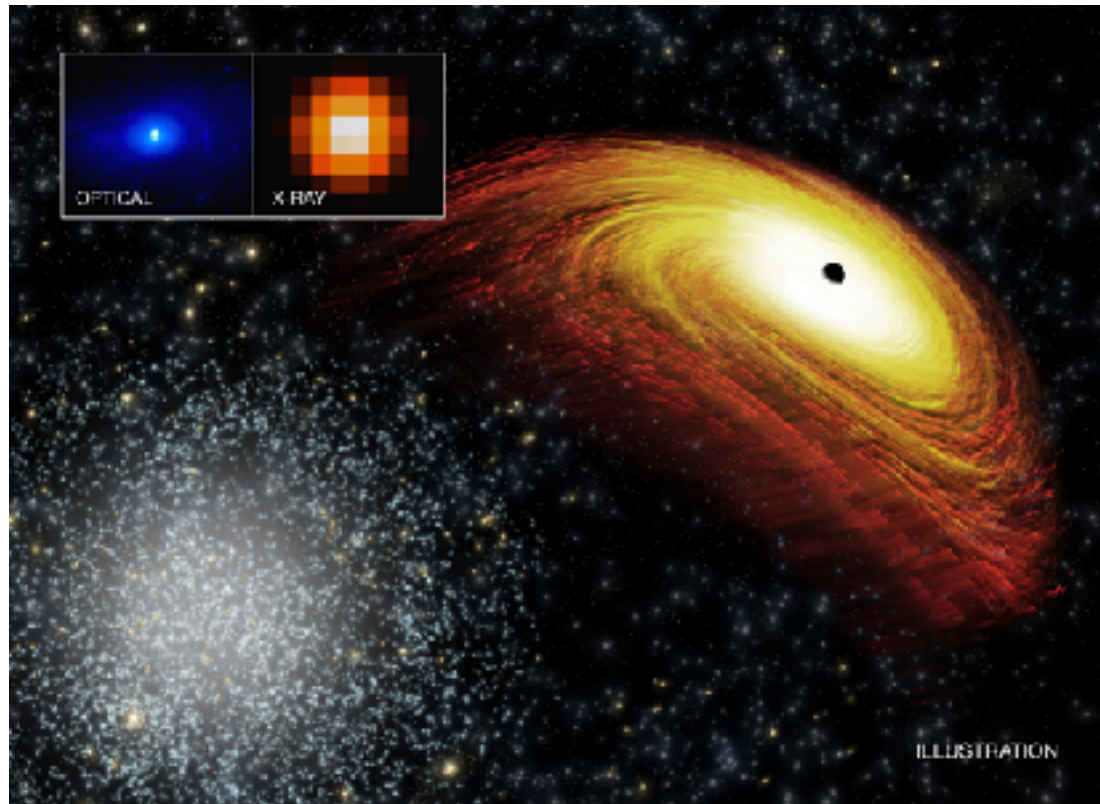
$$v_{\parallel} = K \frac{\eta^2}{(1+q)} \left[ (1 + B_K \eta) |\alpha_2^{\perp} - q \alpha_1^{\perp}| \cos(\Theta_{\Delta} - \Theta_0) \right.$$

$$\left. + K_S \frac{(1-q)}{(1+q)^2} |\alpha_2^{\perp} + q^2 \alpha_1^{\perp}| \cos(\Theta_S - \Theta_1) \right],$$

# Inverse Problem 3: QSO 3C 186 (radio-loud quasar)



Best candidate for GW recoil kicked black hole



11kpc off-set from center ... Broad emission lines  $\rightarrow$   $\sim 2140 \text{ \AA}$  390km/s



## MODELING THE BLACK HOLE MERGER OF QSO 3C 186

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*Draft version April 5, 2017*

### ABSTRACT

Recent detailed observations of the radio-loud quasar 3C 186 indicate the possibility that a supermassive recoiling black hole is moving away from the host galaxy at a speed of nearly 2100km/s. If this is the case, we can model the mass ratio and spins of the progenitor binary black hole using the results of numerical relativity simulations. We find that the black holes in the progenitor must have comparable masses with a mass ratio  $q = m_1/m_2 > 1/4$  and the spin of the primary black hole must be  $\alpha_2 = S_2/m_2^2 > 0.4$ . The final remnant of the merger is bounded by  $\alpha_f > 0.45$  and at least 4% of the total mass of the binary system is radiated into gravitational waves. We consider four different pre-merger scenarios that further narrow those values. Assuming, for instance, a cold accretion driven merger model, we find that the binary had comparable masses with  $q = 0.70_{-0.21}^{+0.29}$  and the normalized spins of the larger and smaller black holes were  $\alpha_2 = 0.94_{-0.22}^{+0.06}$  and  $\alpha_1 = 0.95_{-0.09}^{+0.05}$ . We can also estimate the final recoiling black hole spin  $\alpha_f = 0.93_{-0.03}^{+0.02}$  and that the system radiated 9.6 $_{-1.4}^{+0.8}$ % of its total mass, making the merger of those black holes the most energetic event ever observed.

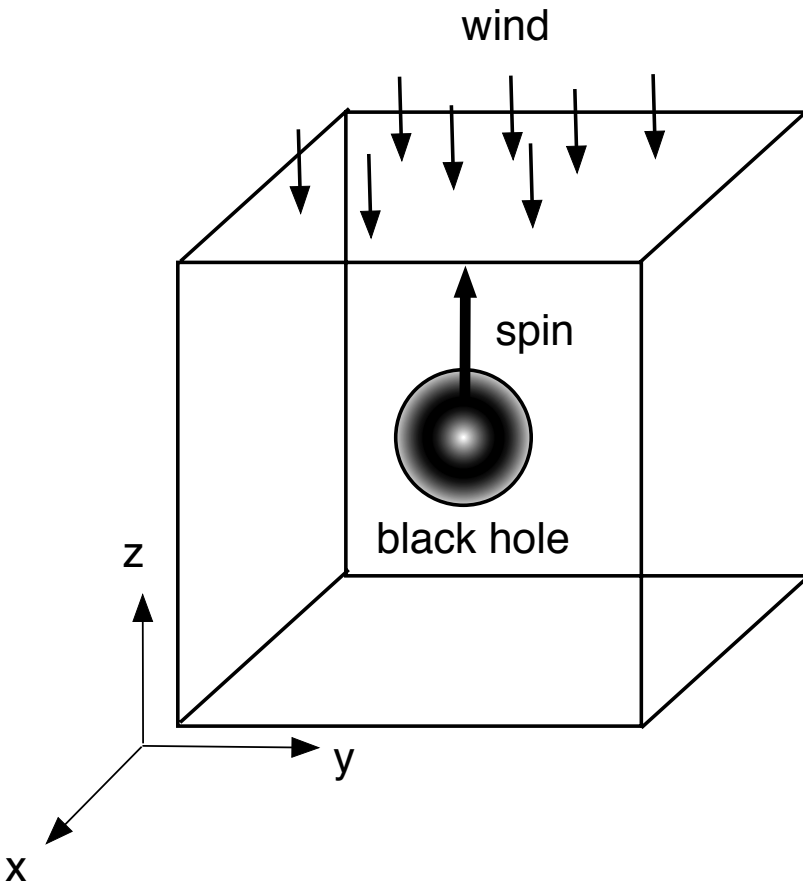
*Keywords:* supermassive black holes — binary merger — gravitational recoils

# Direct problem

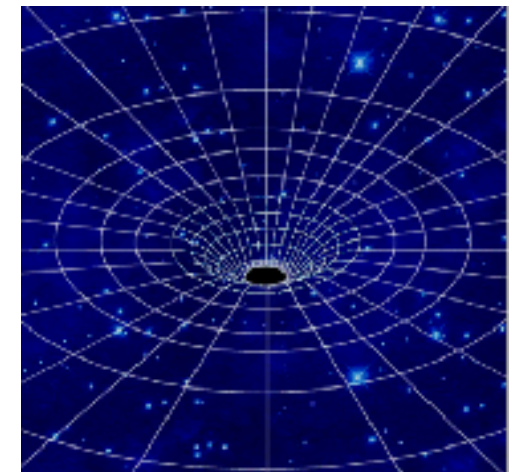


$$T^{\mu\nu} = \rho_0 h u^\mu u^\nu + p g^{\mu\nu}$$

$$\frac{\partial u}{\partial t} + \frac{\partial F^i(u)}{\partial x^i} = S$$



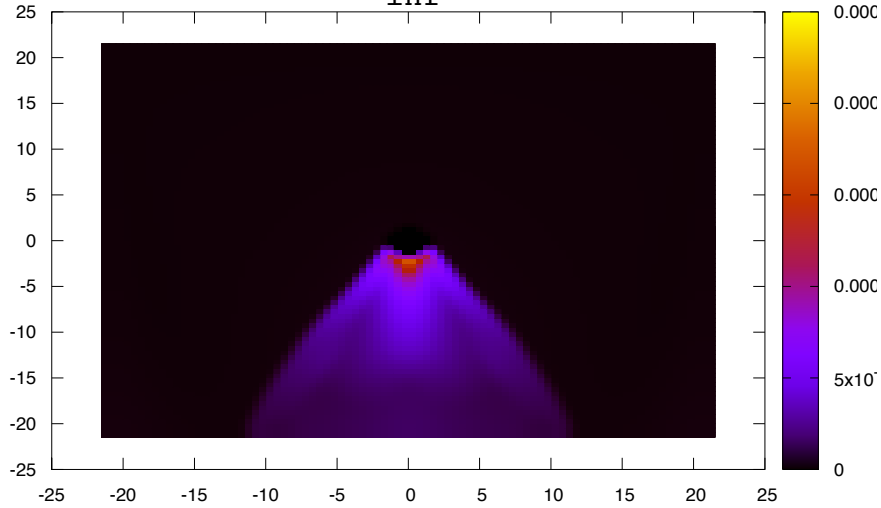
$$u = \begin{bmatrix} D \\ J_1 \\ J_2 \\ J_3 \\ \tau \end{bmatrix}, F^i = \begin{bmatrix} \alpha \left( v^i - \frac{\beta^i}{\alpha} \right) D \\ \alpha \left( v^i - \frac{\beta^i}{\alpha} \right) J_1 + \alpha \sqrt{\gamma} p \delta^i_1 \\ \alpha \left( v^i - \frac{\beta^i}{\alpha} \right) J_2 + \alpha \sqrt{\gamma} p \delta^i_2 \\ \alpha \left( v^i - \frac{\beta^i}{\alpha} \right) J_3 + \alpha \sqrt{\gamma} p \delta^i_3 \\ \alpha \left( v^i - \frac{\beta^i}{\alpha} \right) + \alpha \sqrt{\gamma} v^i p \end{bmatrix}, S = \begin{bmatrix} 0 \\ \alpha \sqrt{\gamma} T^{\mu\nu} g_{\nu\sigma} \Gamma^\sigma_{\mu 1} \\ \alpha \sqrt{\gamma} T^{\mu\nu} g_{\nu\sigma} \Gamma^\sigma_{\mu 2} \\ \alpha \sqrt{\gamma} T^{\mu\nu} g_{\nu\sigma} \Gamma^\sigma_{\mu 3} \\ \alpha \sqrt{\gamma} (T^{\mu 0} \partial_\mu \alpha - \alpha T^{\mu\nu} \Gamma^0_{\mu\nu}) \end{bmatrix}$$



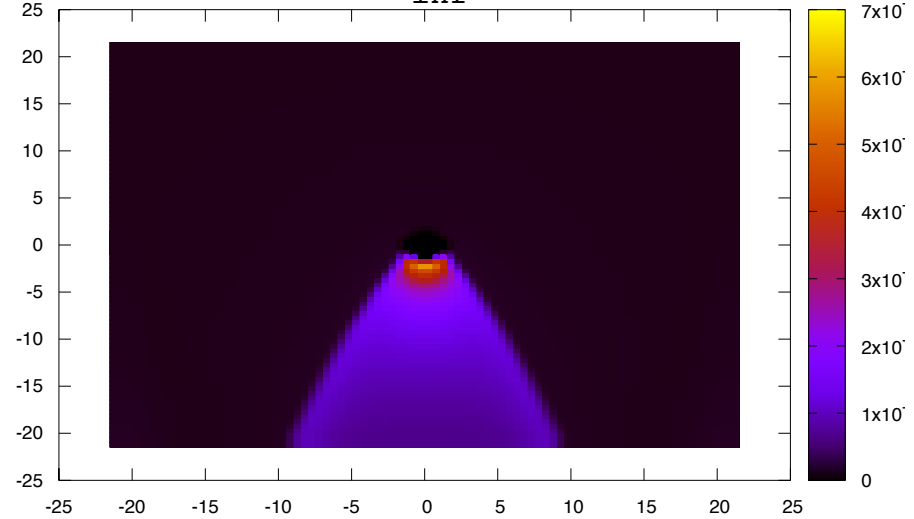
# A sample: 4 out of 900



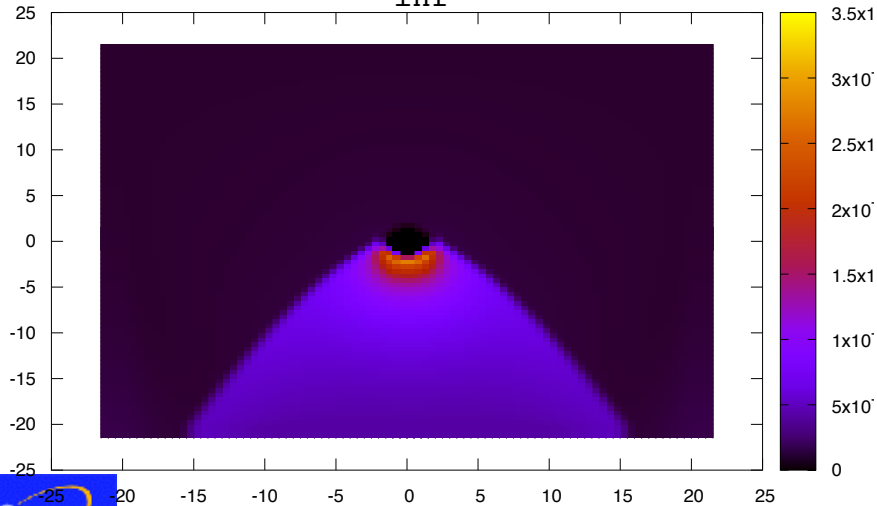
$$\Gamma=1.1 \quad v_{\text{inf}}=0.4c$$



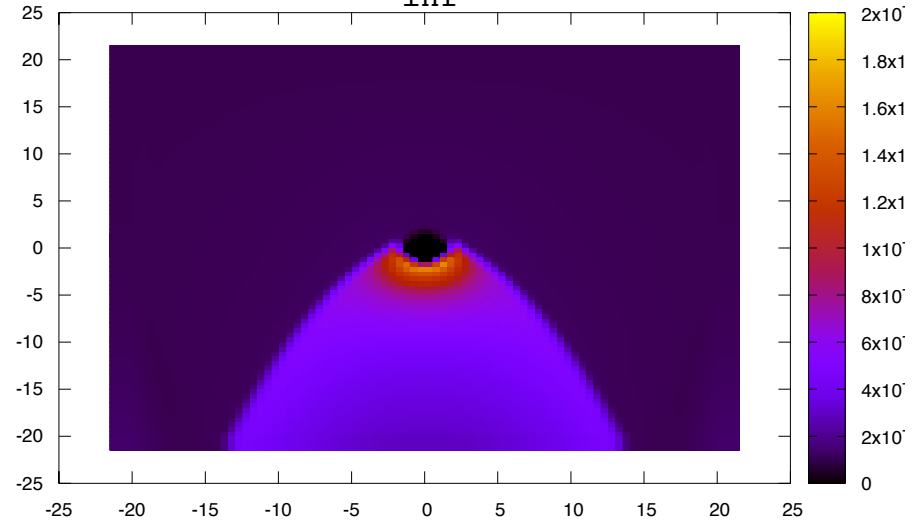
$$\Gamma=1.28 \quad v_{\text{inf}}=0.6c$$



$$\Gamma=1.47 \quad v_{\text{inf}}=0.5c$$



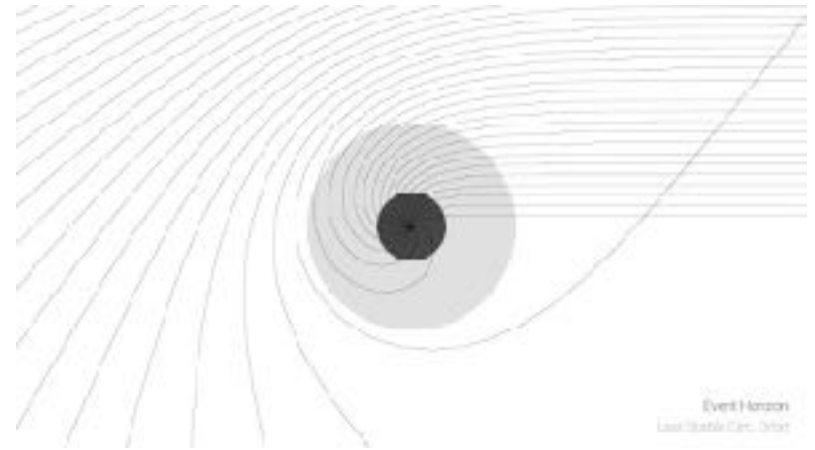
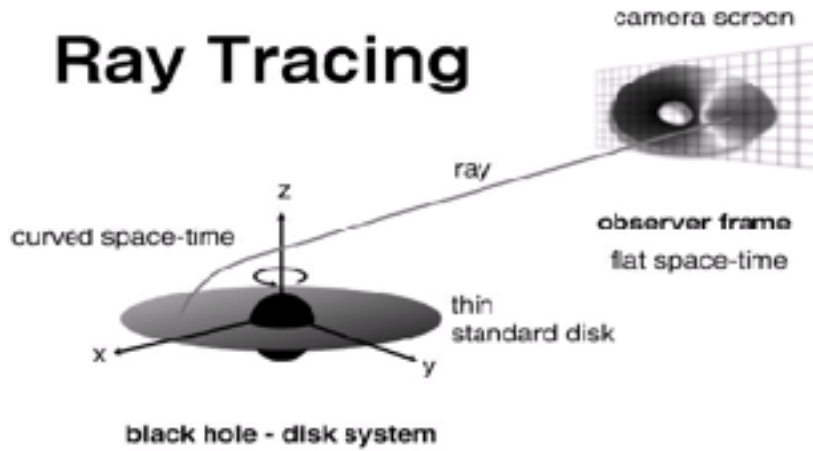
$$\Gamma=5/3 \quad v_{\text{inf}}=0.8c$$



# Ray tracing process

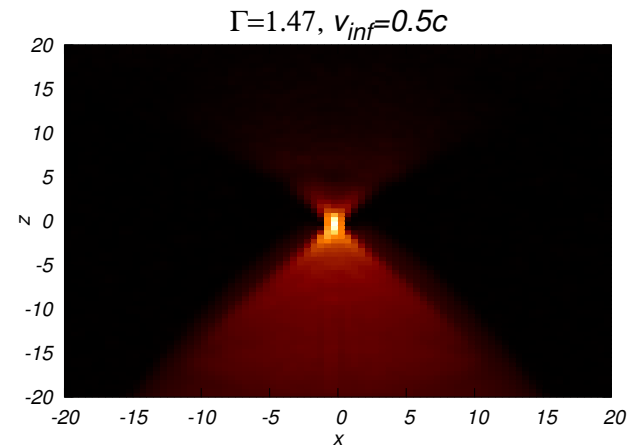
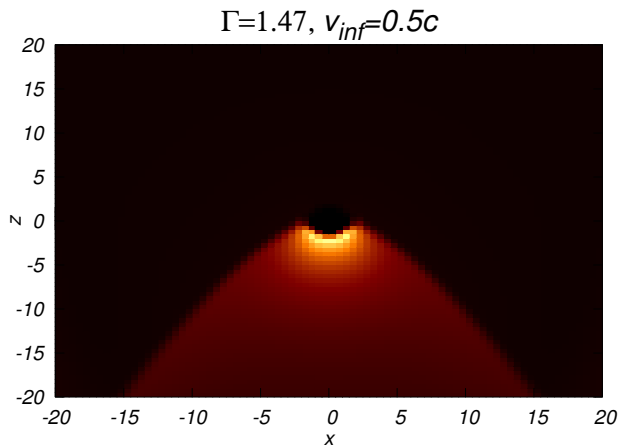
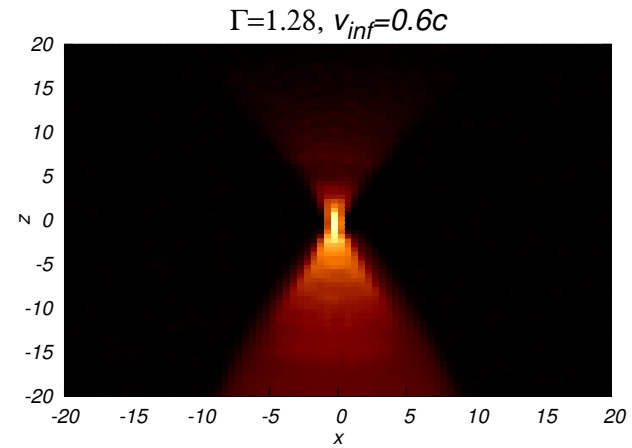
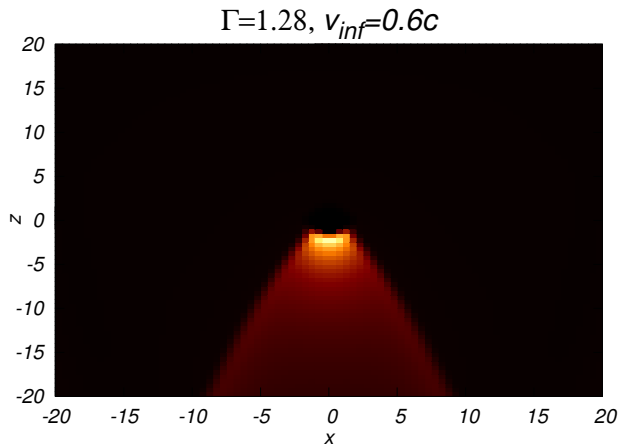


## Ray Tracing





# A little sample 2 out of 900

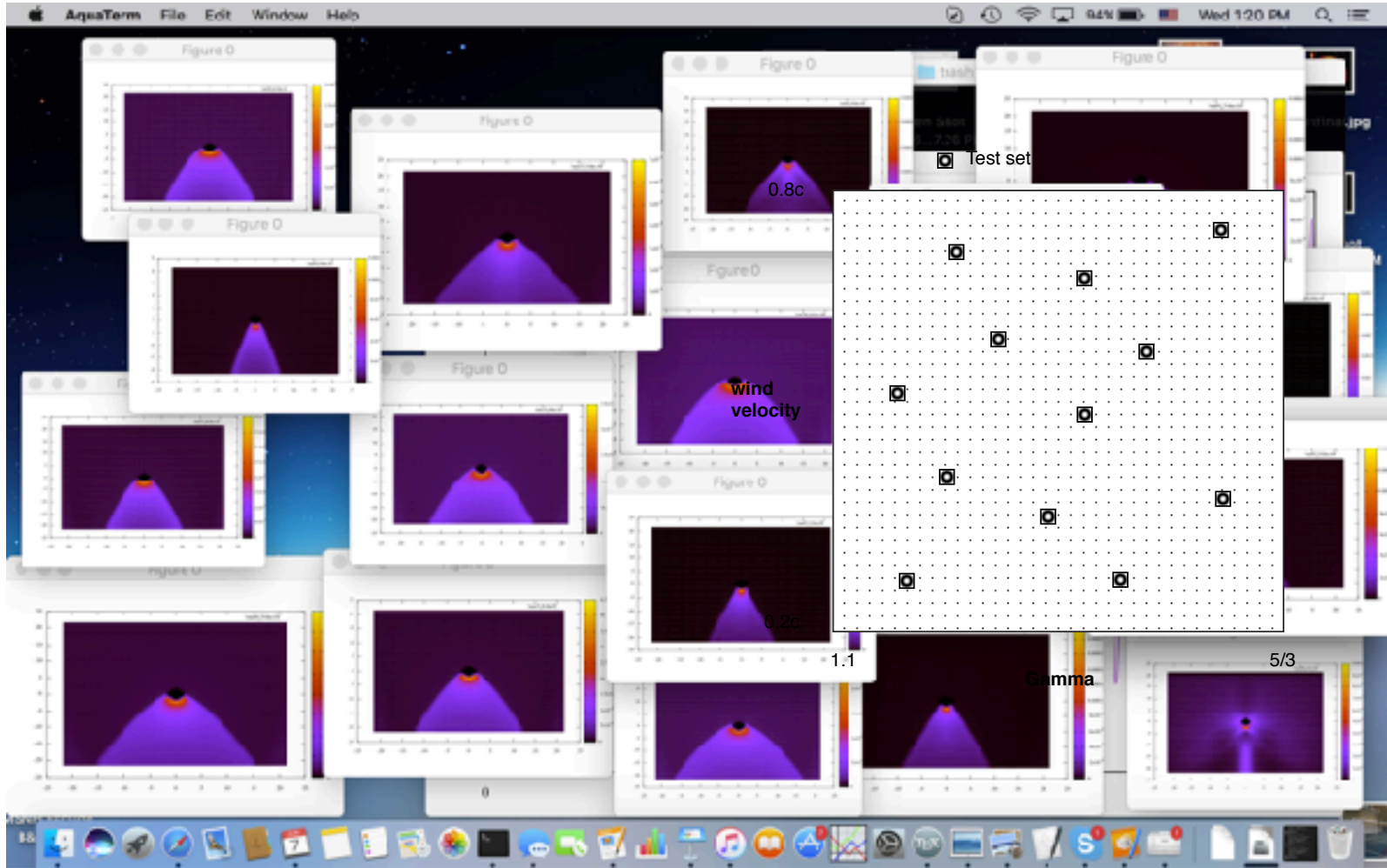


Characterizing the velocity of a wandering black hole and properties of the surrounding medium using convolutional neural networks

J. A González, F. S. Guzmán

Phys. Rev. D 97, 063001 (2018)

# We prepared a sample of 900 of these runs



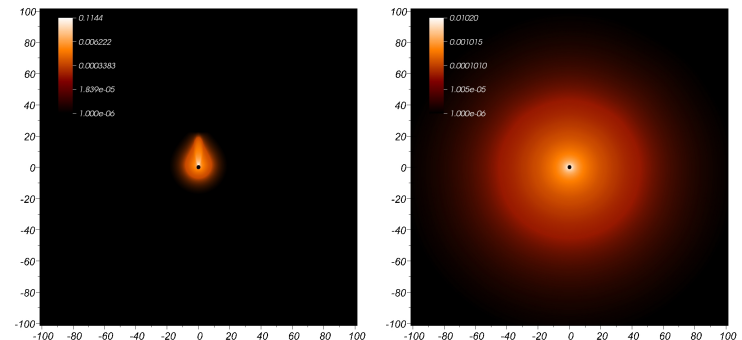
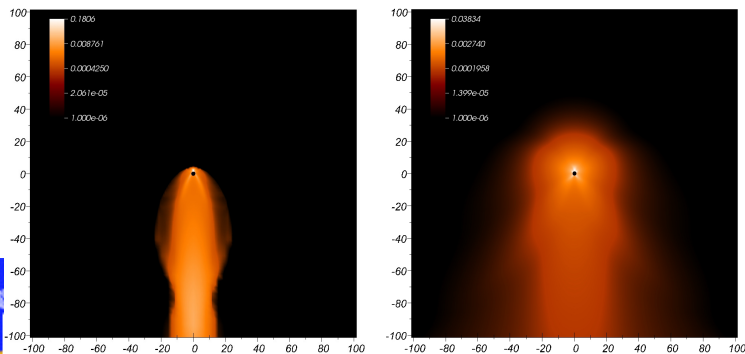
# GOAL



Given an image track the following:

- Properties of the black hole candidates (spin, mass)
- Properties of the matter around (which are model dependent)
- These include equation of state -> degree of ionization
- Temperature - feedbacks the scattering properties
- Opacities (both thermal and scattering)
- Magnetic fields

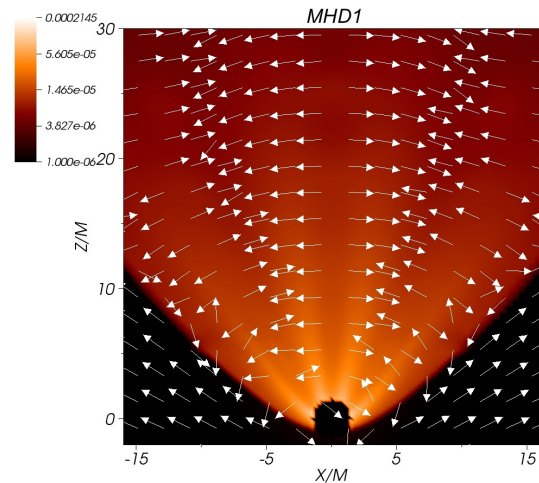
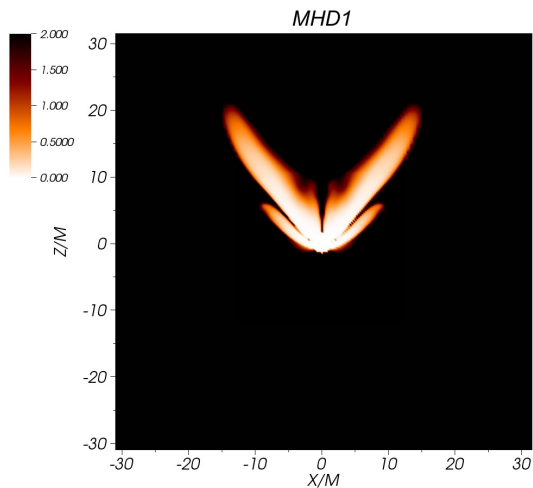
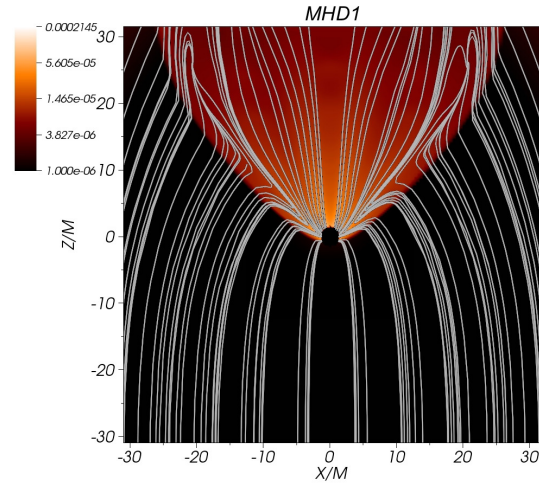
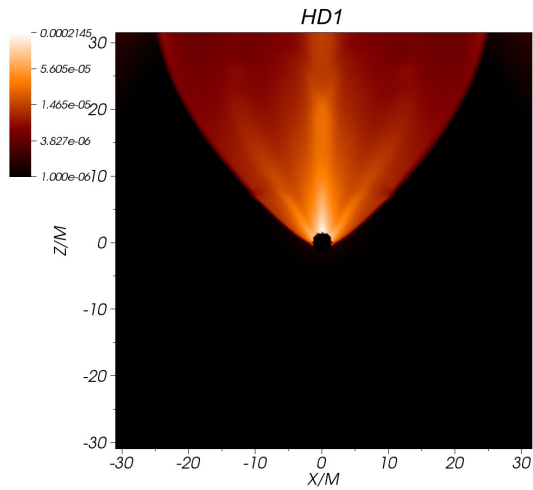
## QSO 3C 186



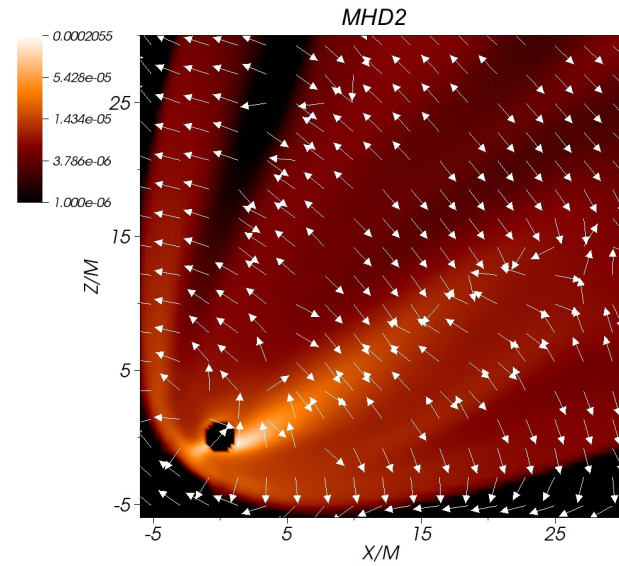
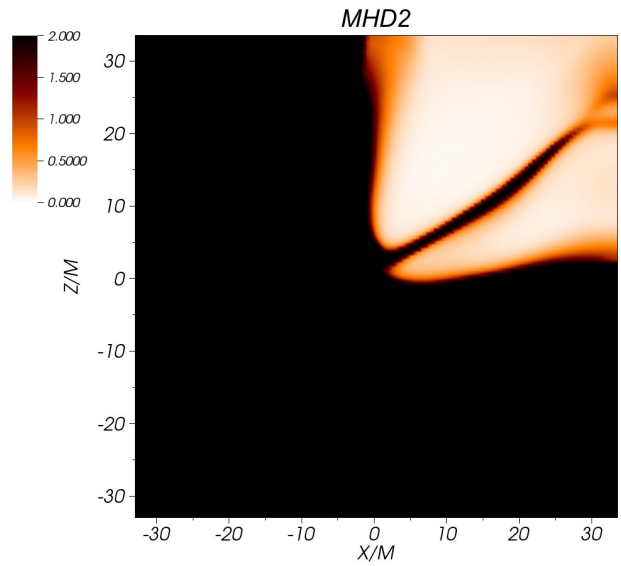
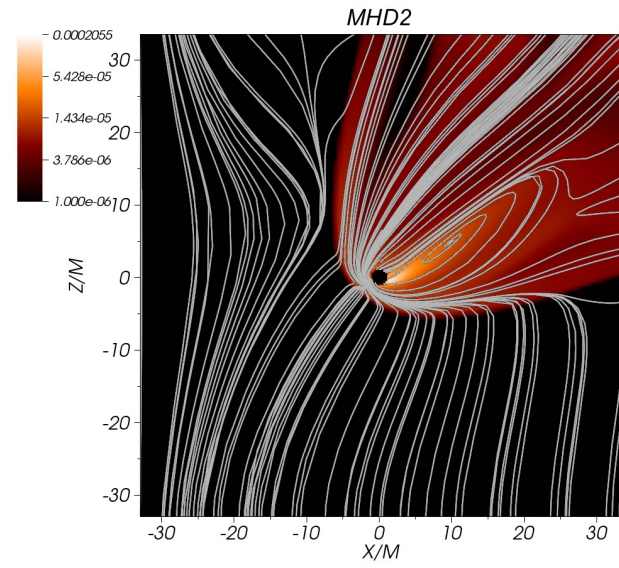
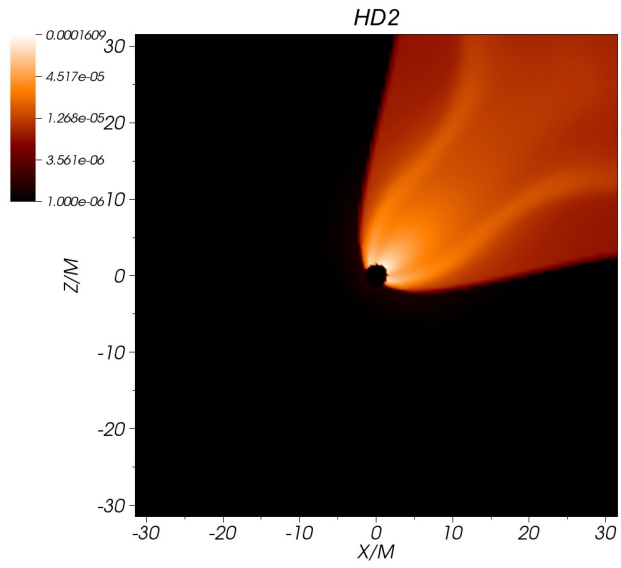
# Magnetized winds



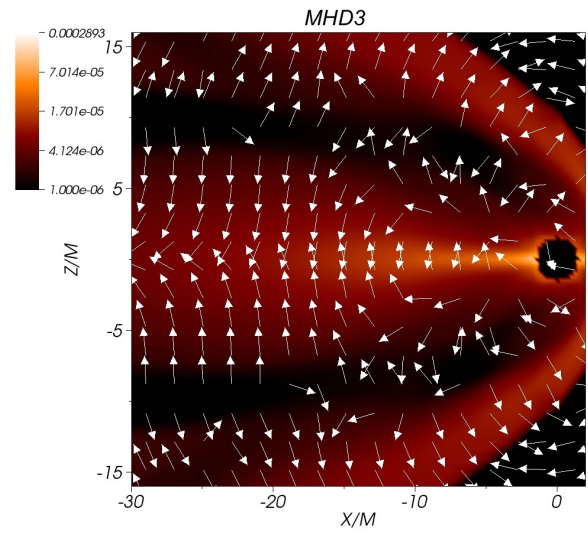
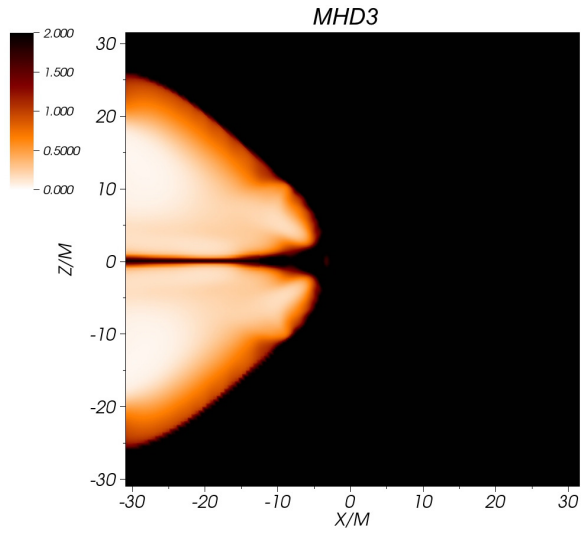
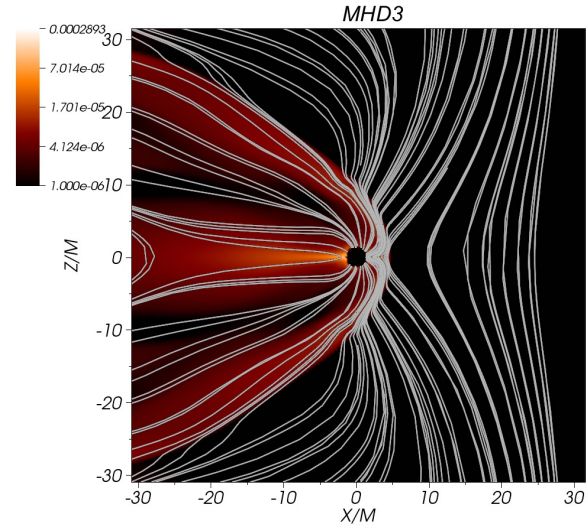
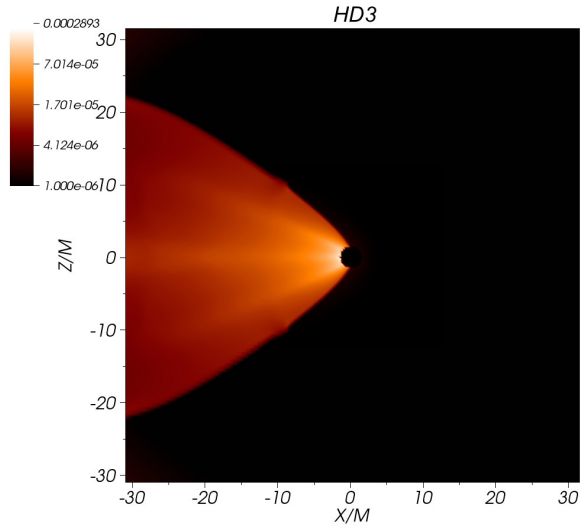
$$T^{\mu\nu} = (\rho h + b^2)u^\mu u^\nu + \left(p + \frac{b^2}{2}\right)g^{\mu\nu} - b^\mu b^\nu$$



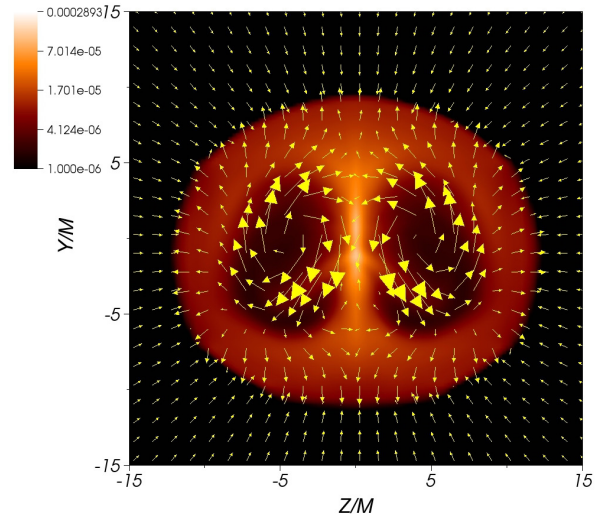
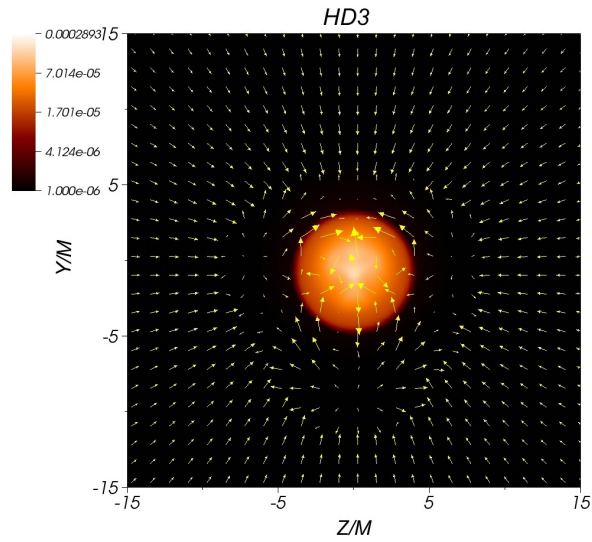
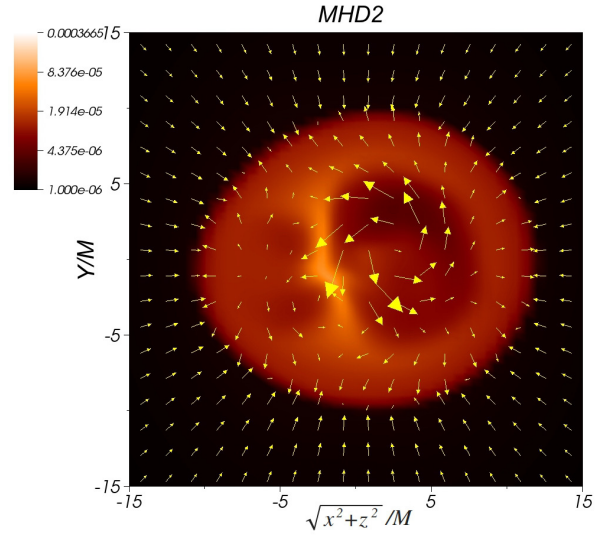
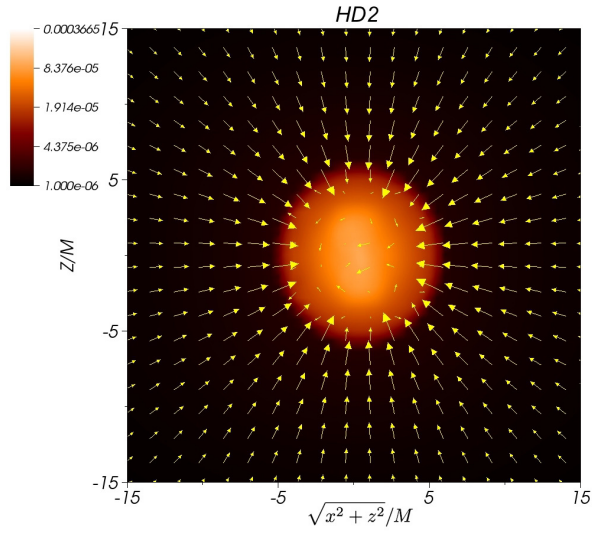
# MHD



# MHD



# MHD



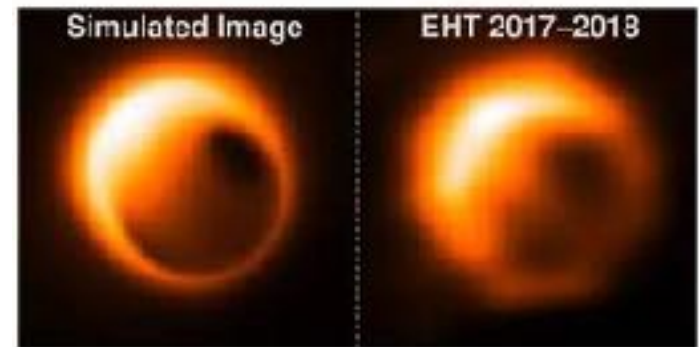
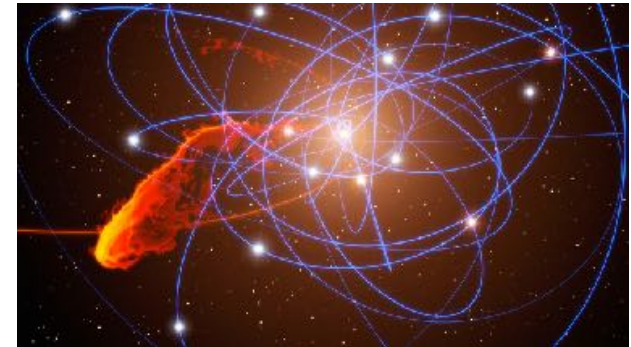
# Inverse Problem 4: EHT Inverse problem



Spin and orientation –  $M$  is estimated from stars around

Is this a black hole?

Is General Relativity ruling there?

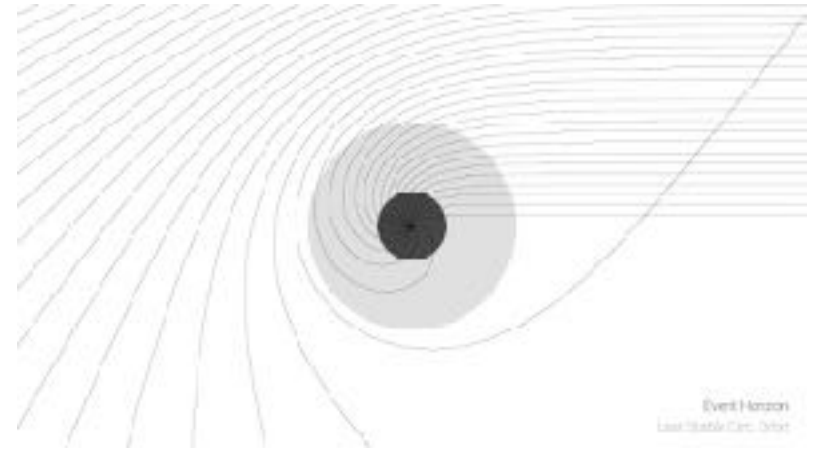
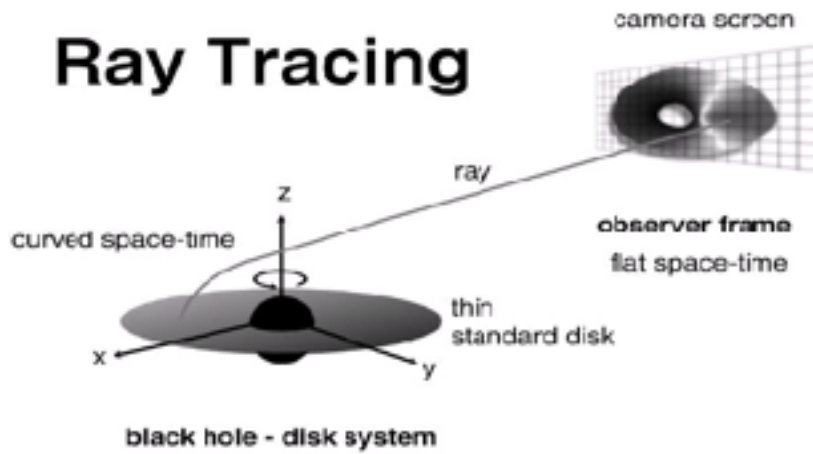




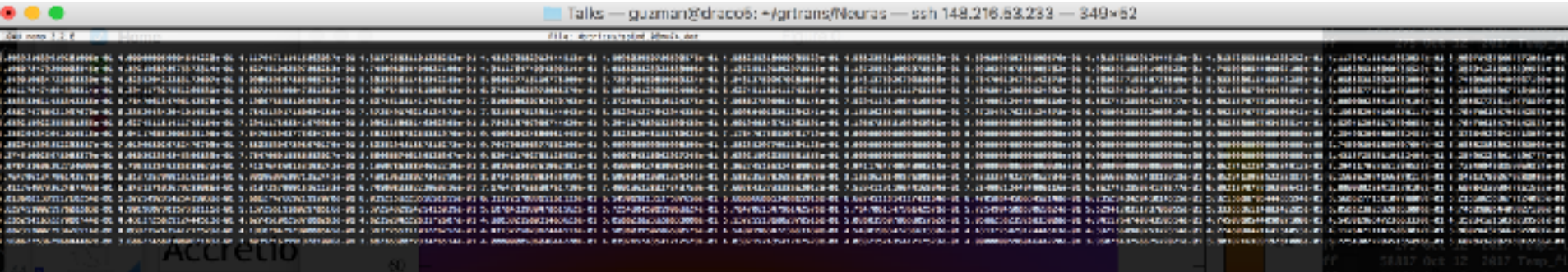
# Ray tracing process



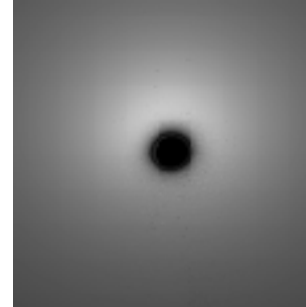
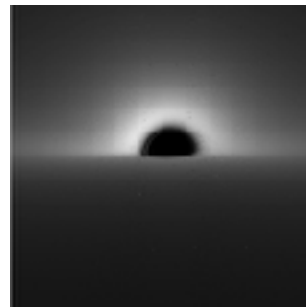
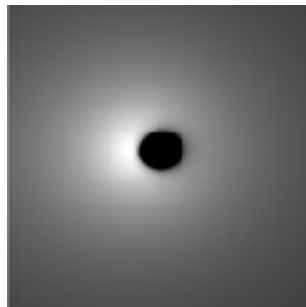
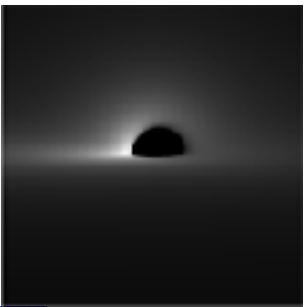
## Ray Tracing



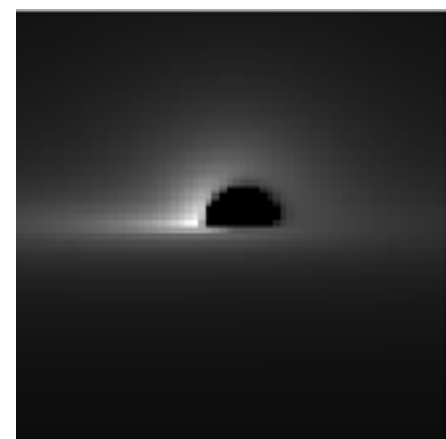
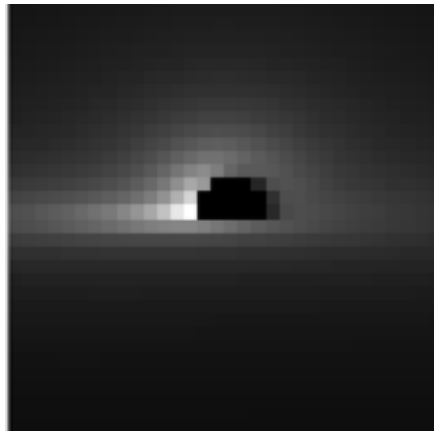
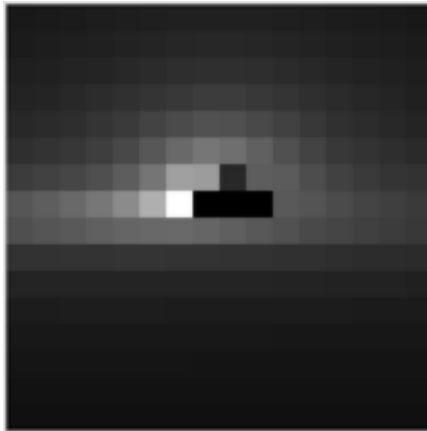
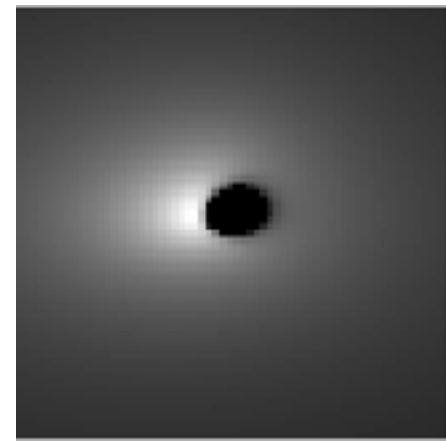
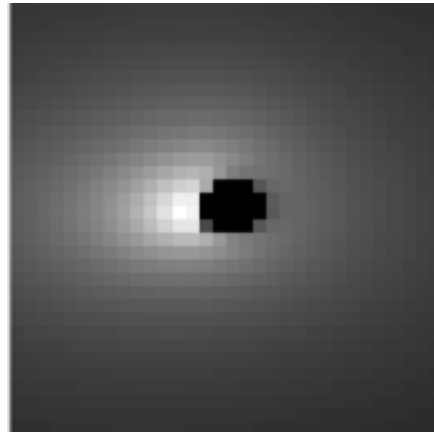
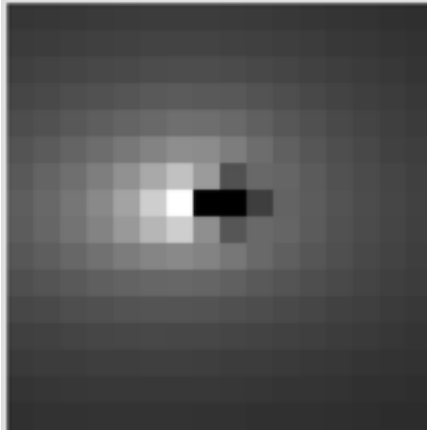
# Accretion



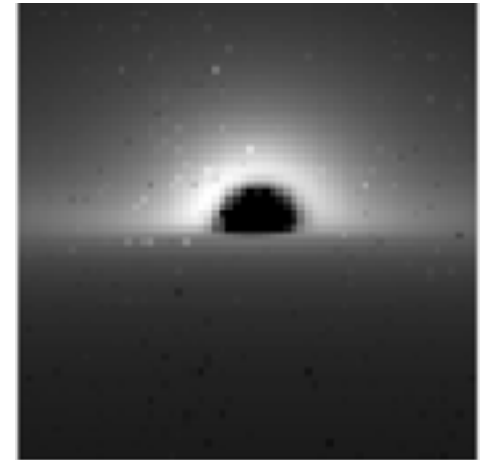
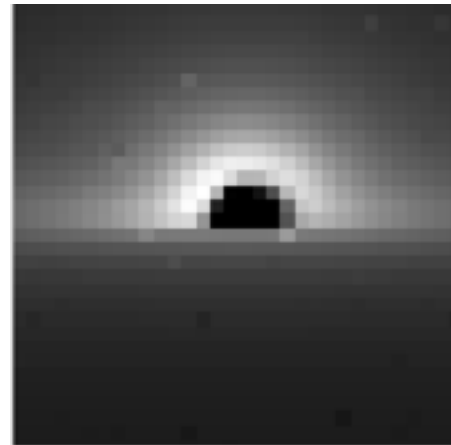
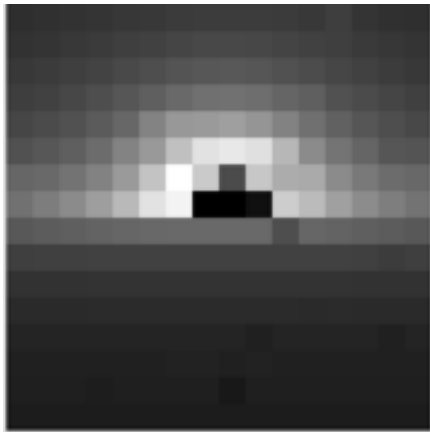
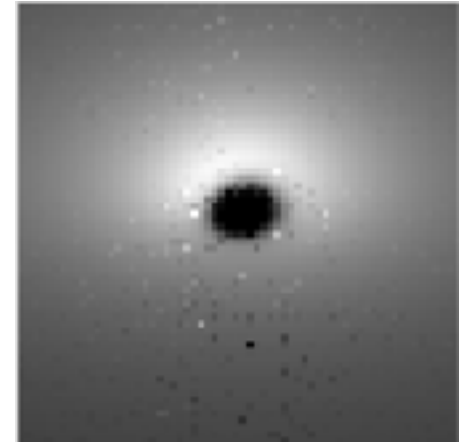
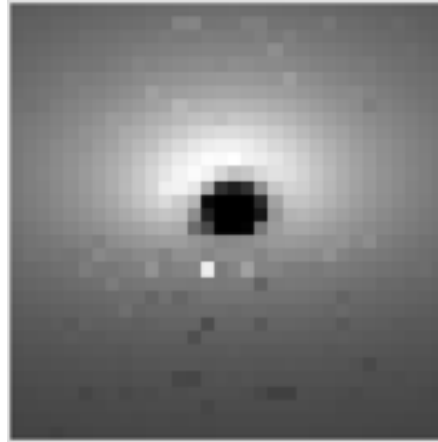
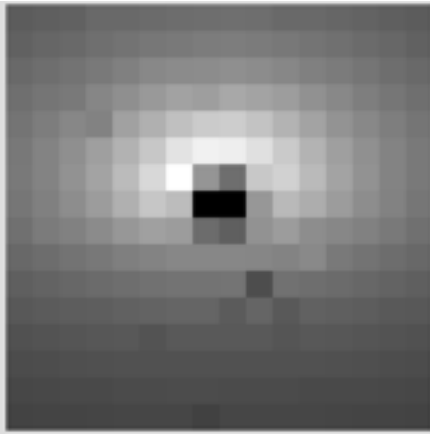
$\nu_{\text{min}} = 7.495e10$  # ← wavelength of 4mm  
 $\nu_{\text{max}} = 3.527e11$  # ← wavelength of 0.85 mm



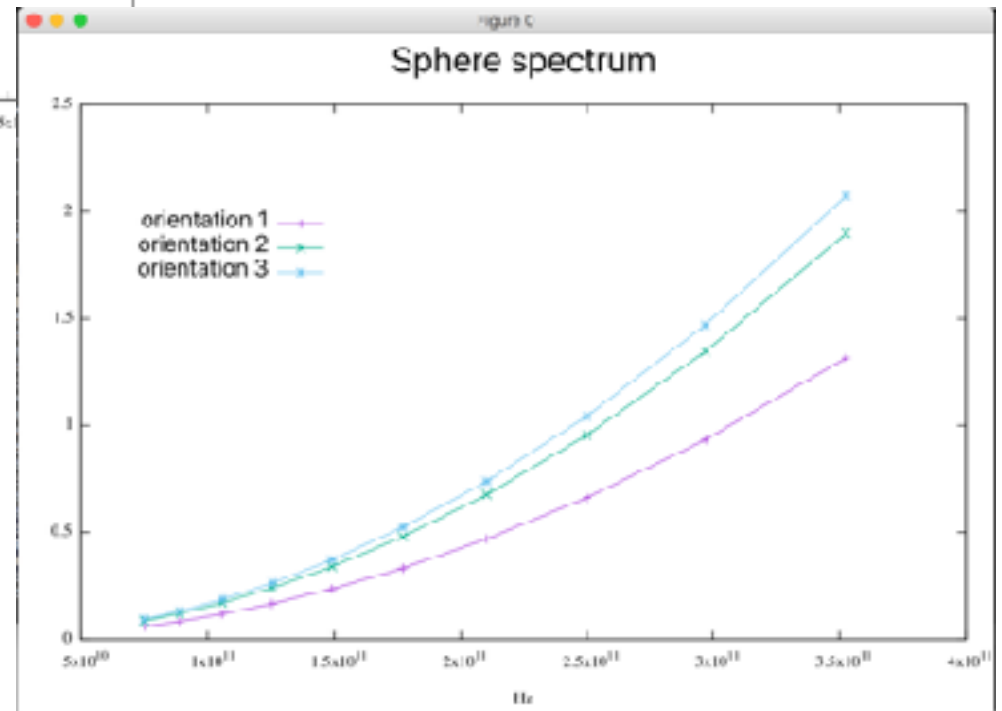
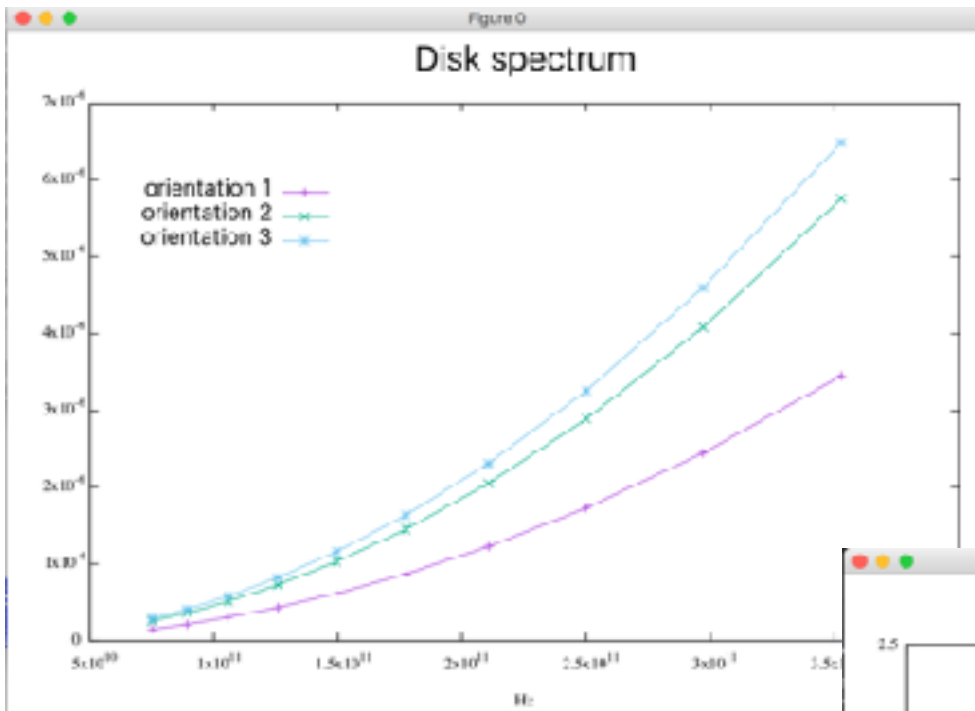
# Accretion disks ( $s=0.85$ , orientation 45, 2)



# Spherical distributions



# Power spectrum, $s=0.95$





# Final comments

In order to solve an IP, one NEEDS to solve DIRECT PROBLEMS

Binary black holes have good data banks generated with PPN approximations

NS-BH Need the solution of the GRMHD equations (very expensive)

The addition of electromagnetic and neutrino radiation to GW sources is the **multi-messenger astronomy, very expensive too... state of the art**

Wandering black holes: simulations and observations needed

Shadows of black holes does not need that much power: observational technology is needed

**WHAT TO DO WITH SOME OF THIS TECHNOLOGY?**

# Template

